

回路設計のためのコンパクトモデルの基本と 先端MOSFETモデルへの拡張

東京工業大学
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広島大学先端物質科学研究科
広島大学HiSIM研究センター
三浦道子

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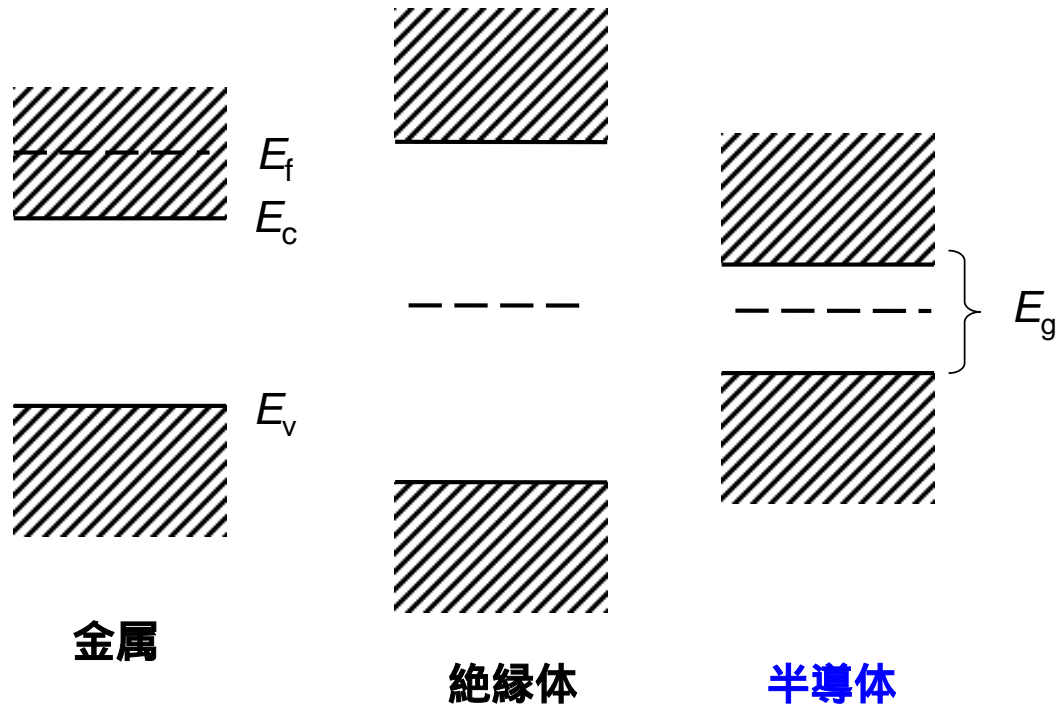
解析式の導出

応用編：

先端デバイスの特性

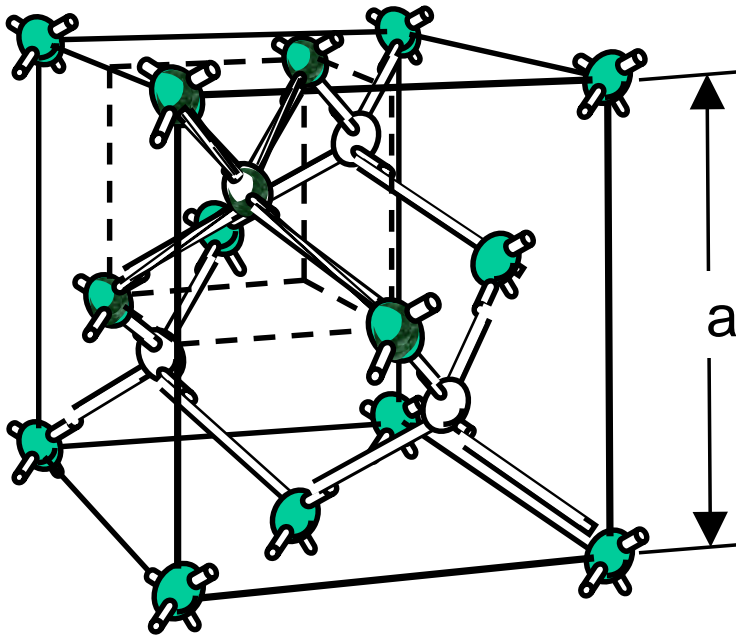
- 短チャネル効果
- Non-Quasi-Static効果
- ノイズ
- 高調波ひずみ
- Y-パラメタ特性

Wilsonのモデル



Bandgap: E_g
Fermi Level: E_f

半導体結晶



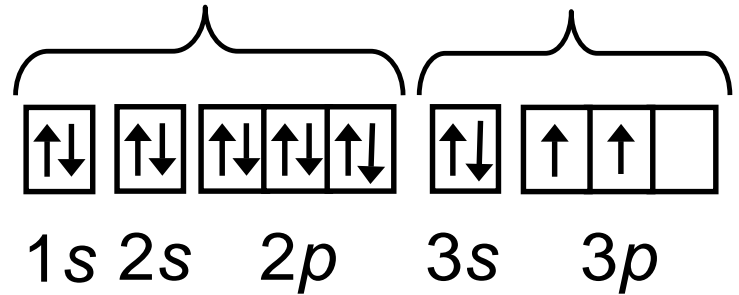
共有結合

電子配置図

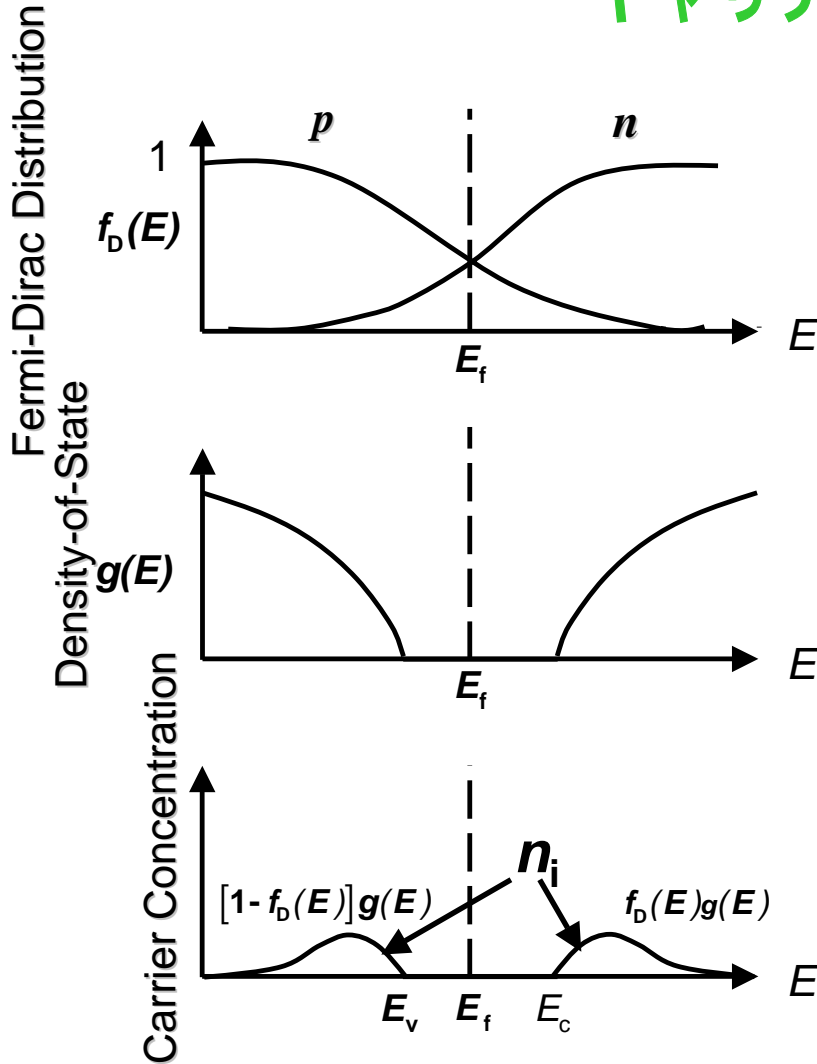
内殻電子

価電子

Si :



キャリア濃度



$$n = \int_{E_c}^{E_{\text{top}}} g(E) f_D(E) dE$$

$$p = \int_{E_{\text{bottom}}}^{E_v} g(E) (1 - f_D(E)) dE$$

$$n = N_c \exp\left(\frac{E_c - E_{fn}}{kT}\right)$$

$$p = N_v \exp\left(\frac{E_{fp} - E_v}{kT}\right)$$

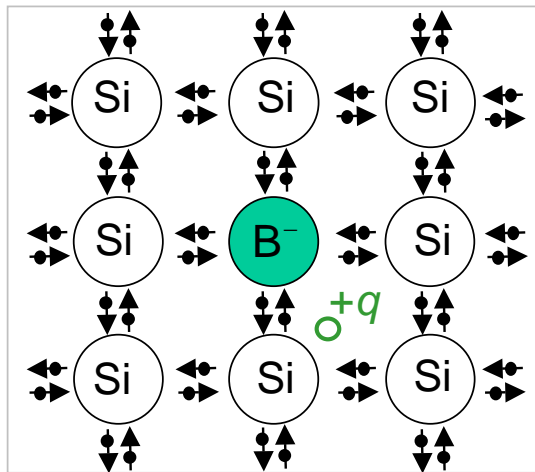
N_c : effective density of state

N_v : “

$$n_i (\text{真性キャリア濃度}) = 1.4 \times 10^{10} / \text{cm}^3$$

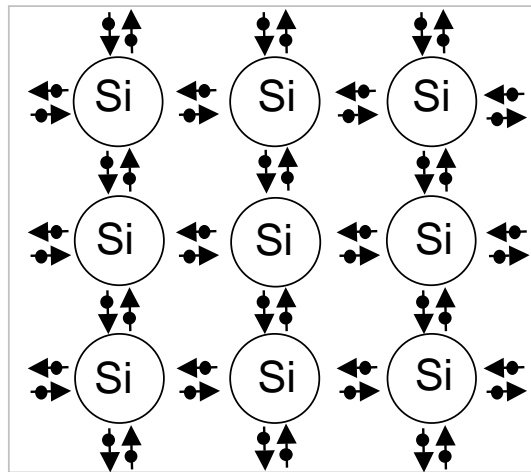
デバイス材料としての半導体

p型シリコン



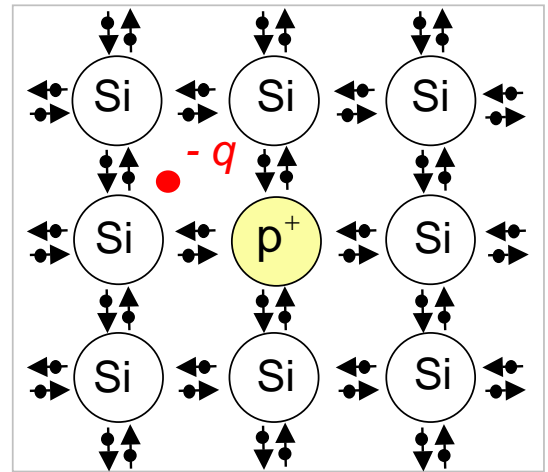
+q: ホール

真性シリコン

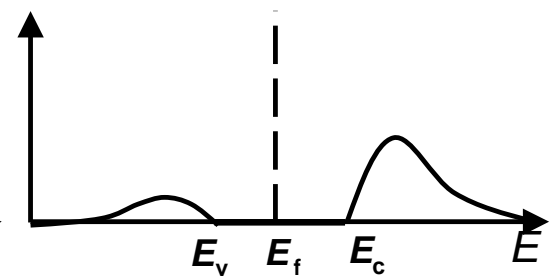
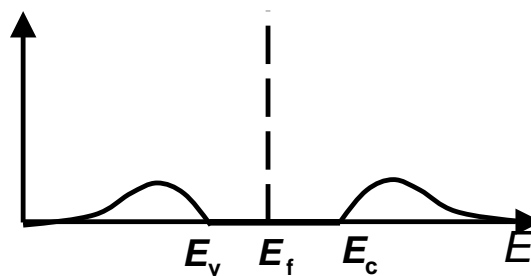
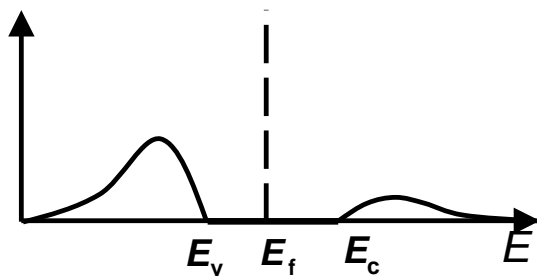


$$n_i (\text{真性キャリア濃度}) = 1.4 \times 10^{10} / \text{cm}^3$$

n型シリコン



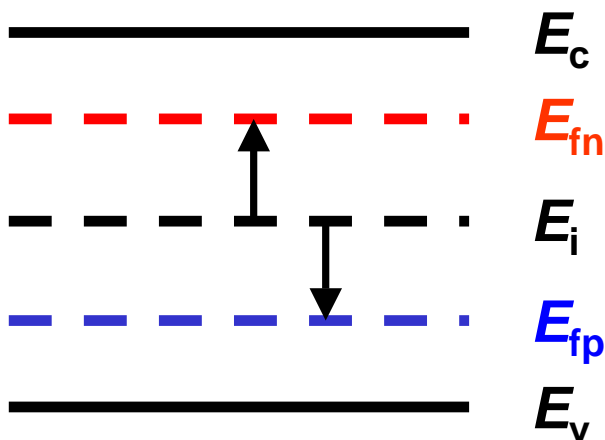
-q: 電子



Fermi順位が変化

$$n = N_c \exp\left(-\frac{E_c - E_{fn}}{kT}\right) = n_i \exp(-q(\phi_{fn} - \phi_i)/kT)$$

$$p = N_v \exp\left(-\frac{E_{fp} - E_v}{kT}\right) = n_i \exp(-q(\phi_i - \phi_{fp})/kT)$$



$$E = -q\phi$$

$$N_c N_v \exp(E_g/kT) = n_i^2$$

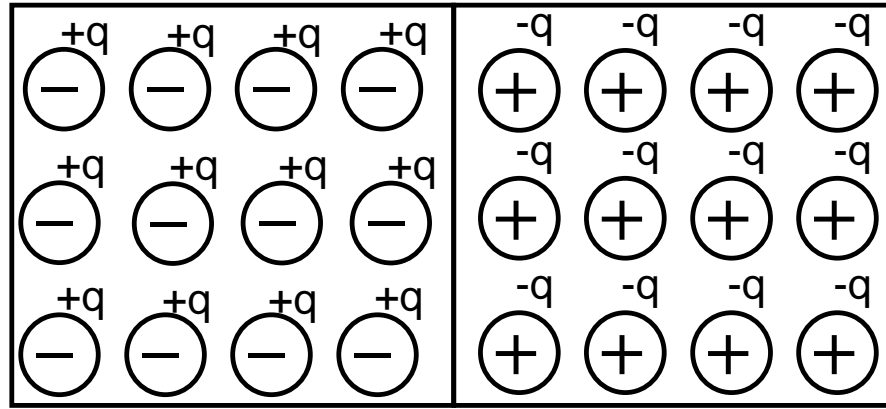
平衡状態 : $pn = n_i^2$

p/n接合

p型シリコン

n型シリコン

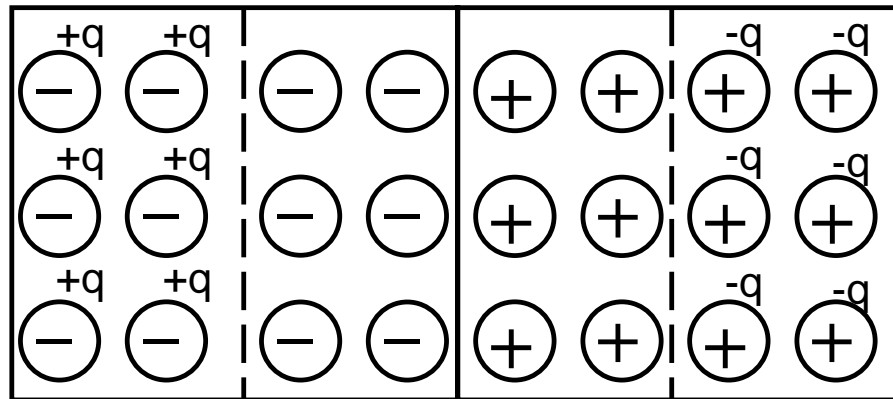
p/n Junction



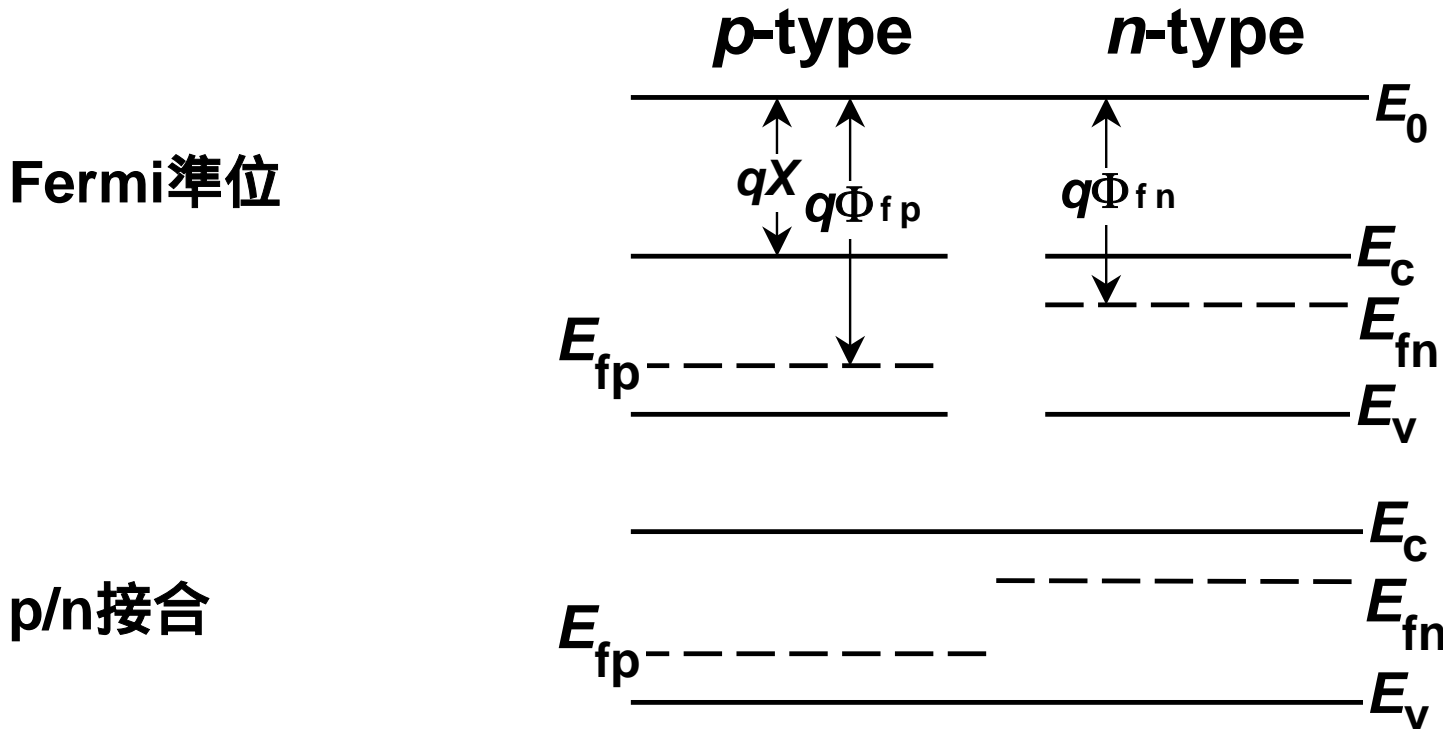
⊖ : N_A
⊕ : N_D

p/n接合

Equilibrium Condition

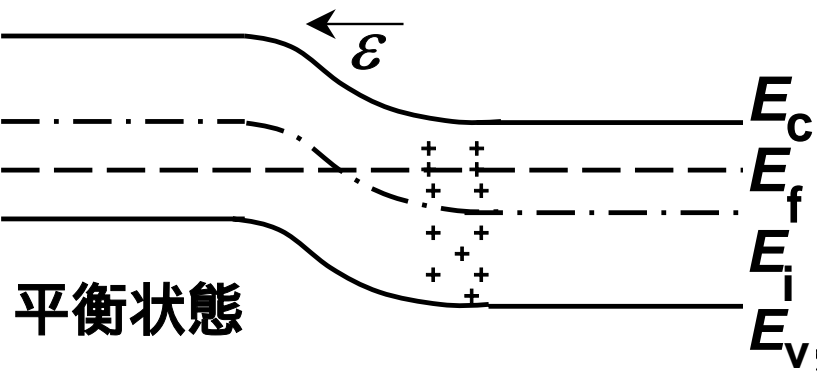


空乏層 (depletion region)

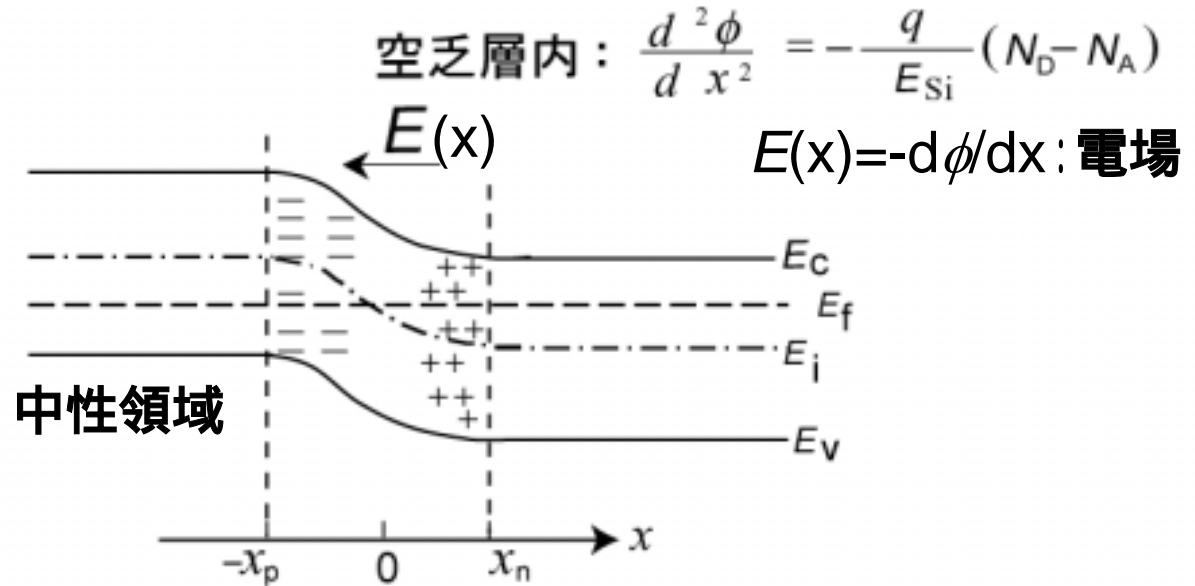


空乏層 (Depletion Region) の形成

- 電場の強さ
- 空乏層幅



Poisson方程式



(i) $0 \leq x \leq x_n$

$$E(x) = -\frac{q N_D}{\epsilon_{Si}} (x_n - x) ; \phi(x) = \phi_n - \frac{q N_D}{2 \epsilon_{Si}} (x_n - x)^2$$

(ii) $-x_p \leq x \leq 0$

$$E(x) = -\frac{q N_A}{\epsilon_{Si}} (x + x_p) ; \phi(x) = \phi_p + \frac{q N_A}{2 \epsilon_{Si}} (x + x_p)^2$$

空乏層幅： $x_n + x_p = \left[\frac{2 \epsilon_{Si}}{q} \phi_i \left(\frac{1}{N_D} + \frac{1}{N_A} \right) \right]^{\frac{1}{2}} ; \phi_i = \phi_n - \phi_p$
(Built-in Potential)

デバイスの基本方程式

-Poisson:
$$\nabla^2 \phi = -\frac{q}{\epsilon_S} (N_D - N_A + p - n)$$

$$n = n_i \exp \frac{q(\phi - \phi_n)}{kT}$$

$$p = n_i \exp \frac{q(\phi_p - \phi)}{kT}$$

-Current Density:
$$j_n = -q \mu_n n \frac{\phi}{y} + q D_n \nabla n$$

$$j_p = -q \mu_p p \frac{\phi}{y} - q D_p \nabla p$$

-Continuity:
$$\frac{\partial n}{\partial t} = G_n - R_n + \frac{1}{q} \nabla j_n = 0$$

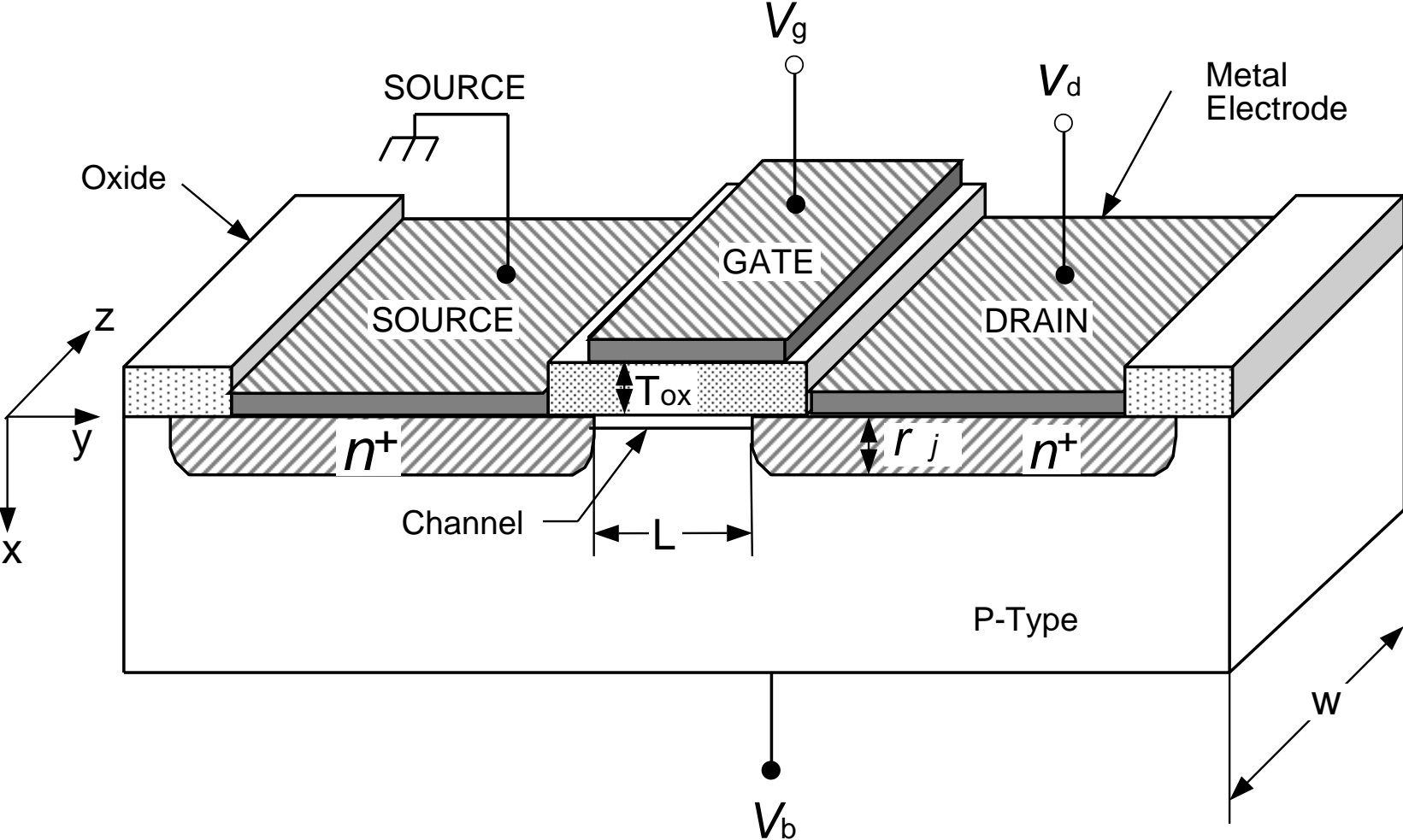
$$\frac{\partial p}{\partial t} = G_p - R_p - \frac{1}{q} \nabla j_p = 0$$

回路モデルとは

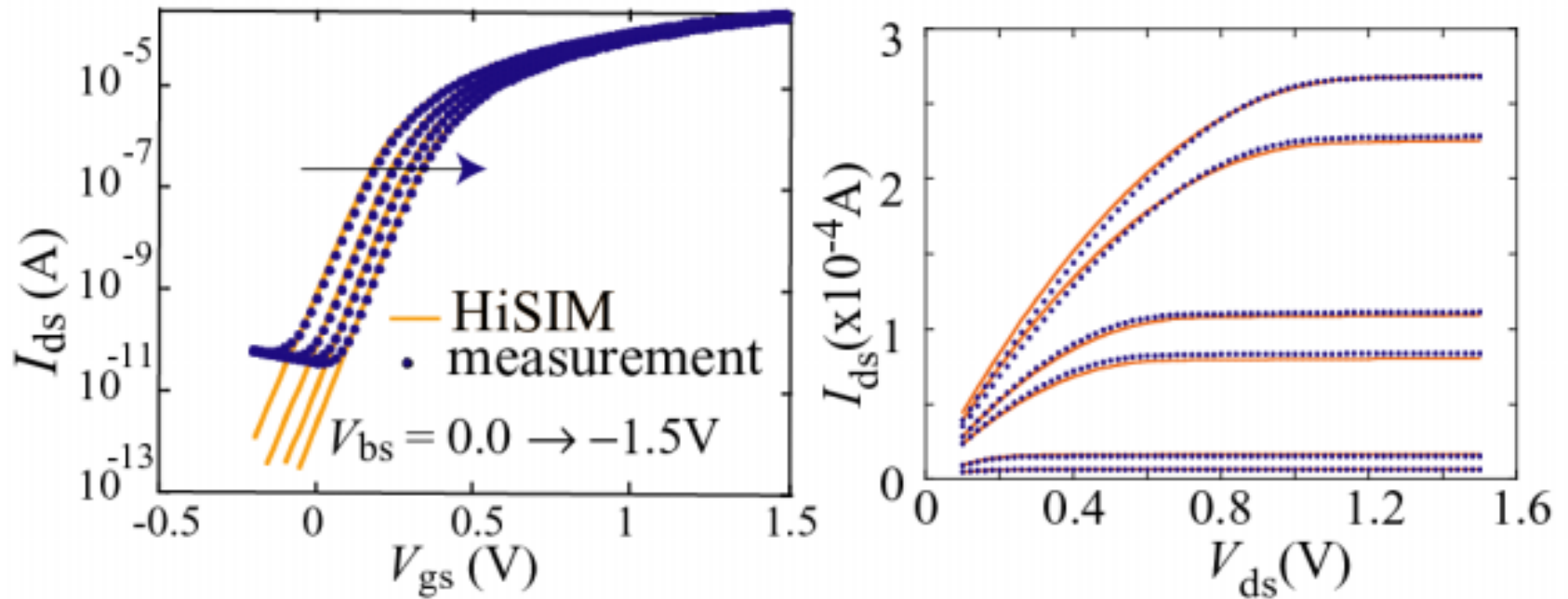
デバイス特性を解析式で記述したもの

- ✓ デバイス特性がコンパクトに記述されている
- ✓ すべての観測される現象がモデル化されている
- ✓ デバイス特性と回路性能の関係が予測できる
- ✓ 大規模回路がシミュレーションできる

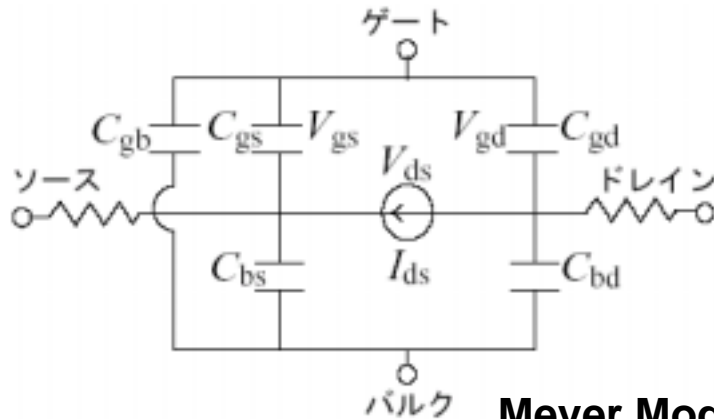
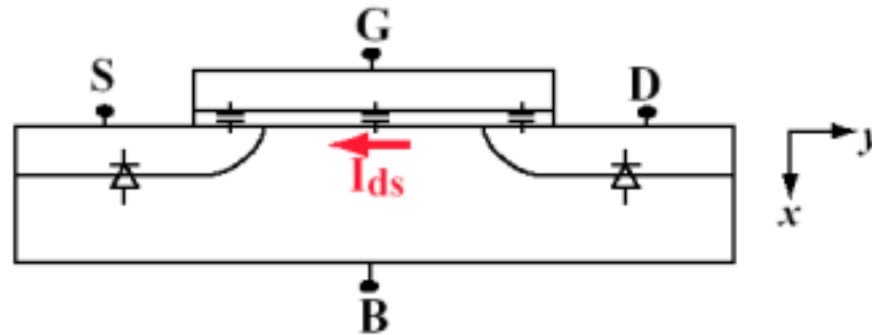
MOSFET



電流の測定値



MOSFET用回路モデル



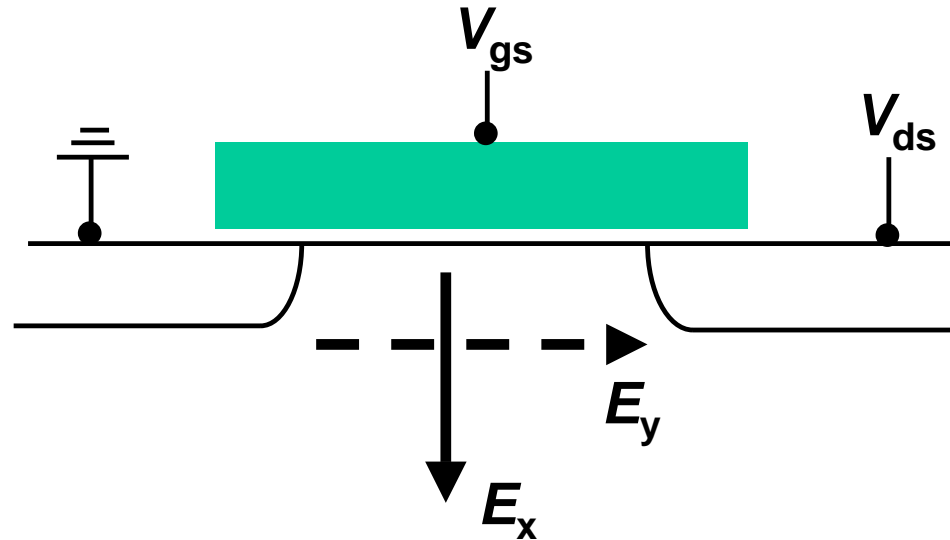
Meyer Model

J. E. Meyer, RCA Rev., vol. 32, p. 42, 1971.

$$I_{ds} = \mu \frac{W}{L} C_{ox} \left[(V_{gs} - V_{th}) V_{ds} - \frac{1}{2} V_{ds}^2 \right] \Rightarrow \text{ドリフト近似}$$

C. T. Sah, IEEE TED, ED11, 324, 1964.

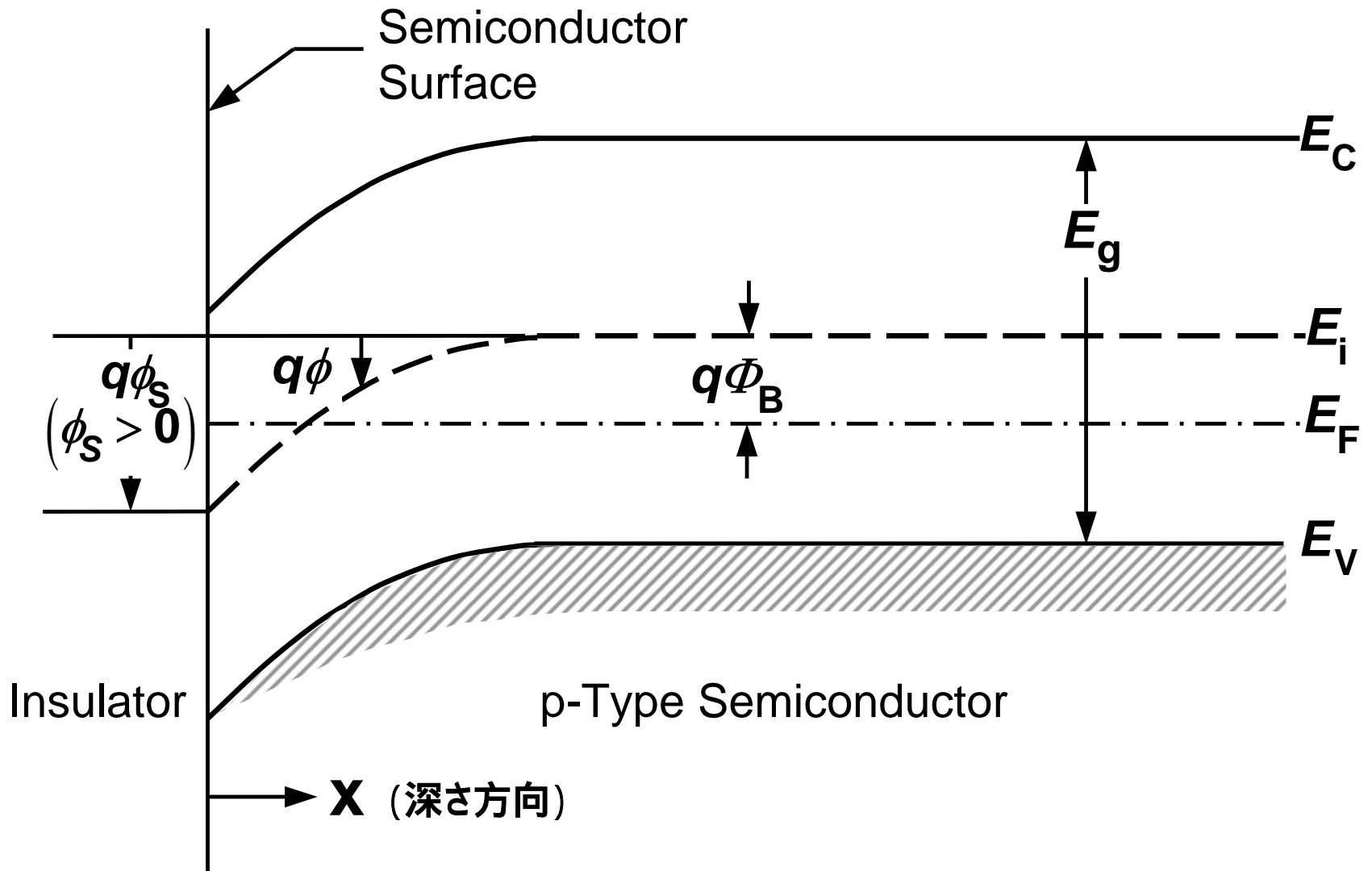
MOSFETの特徴

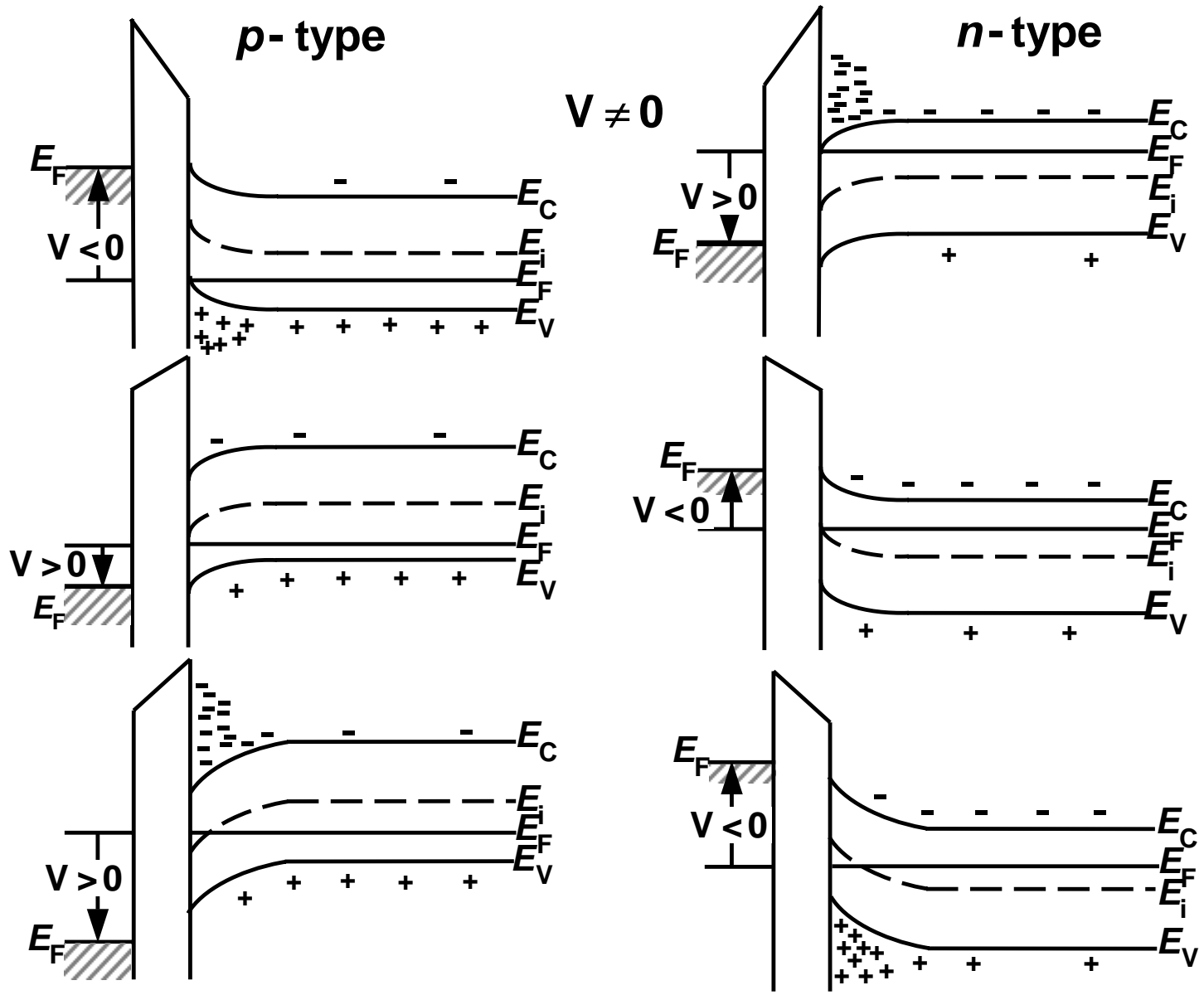


電流: $I = qnv$ $\left\{ \begin{array}{l} q: \text{電荷} \\ n: \text{キャリア濃度} \\ v: \text{キャリア速度} \end{array} \right.$

V_{gs} : キャリアの濃度を定める
 V_{ds} : キャリアの速さを定める

V_{gs}に対するポテンシャル分布





デバイスの基本方程式

-Poisson:
$$\nabla^2 \phi = -\frac{q}{\epsilon_S} (N_D - N_A + p - n)$$

$$n = n_i \exp \frac{q(\phi - \phi_n)}{kT}$$

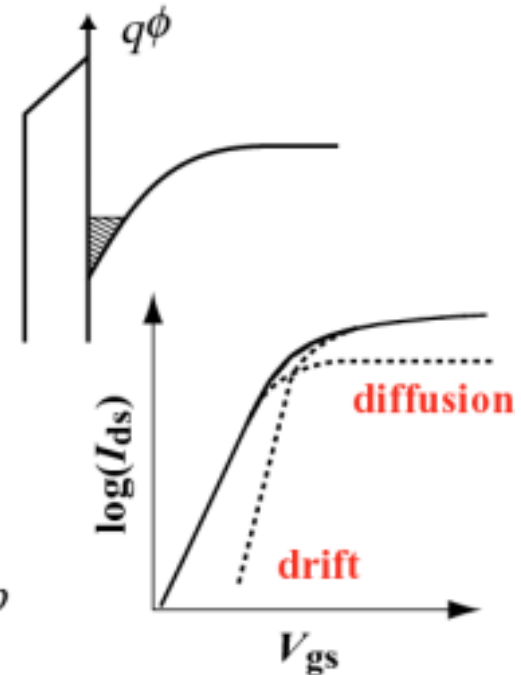
$$p = n_i \exp \frac{q(\phi_p - \phi)}{kT}$$

-Current Density:
$$j_n = -q \mu_n n \frac{\phi}{y} + q D_n \nabla n$$

$$j_p = -q \mu_p p \frac{\phi}{y} - q D_p \nabla p$$

-Continuity:
$$\frac{\partial n}{\partial t} = G_n - R_n + \frac{1}{q} \nabla j_n = 0$$

$$\frac{\partial p}{\partial t} = G_p - R_p - \frac{1}{q} \nabla j_p = 0$$



Poisson方程式: ϕ_s

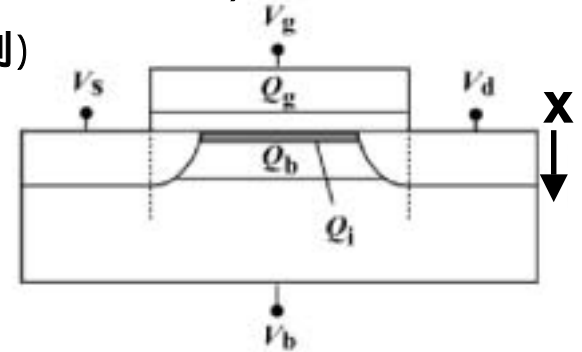
$$\frac{d^2\phi}{dx^2} = -\frac{\rho(x)}{\epsilon_s} \quad (\text{Gradual Channel近似: 深さ方向のみ考慮する})$$

(平衡状態: 電荷中性の法則)

$$\rho(x) = q(N_D - N_A + p - n)$$

$$N_D - N_A = n_{p0} - p_{p0}$$

$$p - n = p_{p0} \exp(-\beta\phi) - n_{p0} \exp(\beta\phi)$$



$$\frac{\partial^2\phi}{\partial x^2} = -\frac{q}{\epsilon_s} \int \left\{ p_{p0} [\exp(-\beta\phi) - 1] - n_{p0} [\exp(\beta\phi) - 1] \right\}$$

$$\int_0^{\partial\phi/\partial x} \left(\frac{\partial\phi}{\partial x} \right) d \left(\frac{\partial\phi}{\partial x} \right) = -\frac{q}{\epsilon_s} \int_0^{\phi} \left\{ p_{p0} [\exp(-\beta\phi) - 1] - n_{p0} [\exp(\beta\phi) - 1] \right\} d\phi$$

$$E_s^2 = \left(\frac{2kT}{q} \right)^2 \left(\frac{qp_{p0}\beta}{2\epsilon_{si}} \right) \left\{ \left[\exp(-\beta\phi) + \beta\phi + 1 \right] + \frac{n_{p0}}{p_{p0}} \left[\exp(\beta\phi) - \beta\phi - 1 \right] \right\}$$

Gauss's Law: $E_{ox} \epsilon_{ox} = E_s \epsilon_s = -Q_s$ Q_s : 基板中の全電荷

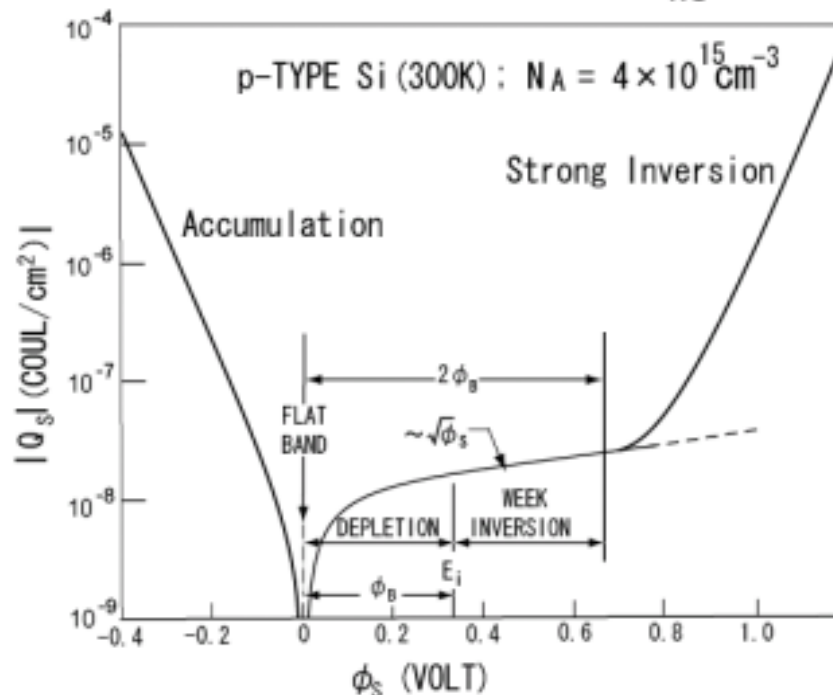
$$E_{ox} = (V_{gs} - V_{fb} - \phi_s) / T_{ox}$$

MOSキャパシタンス

$$-Q_S = C_{ox}(V'_G - \phi_S(y)) = \sqrt{\frac{2\epsilon_s q N_{sub}}{\beta}} \left[\exp\{-\beta(\phi_S(y) - V_{bs})\} + \beta(\phi_S(y) - V_{bs}) - 1 + \frac{n_{p0}}{p_{p0}} \left\{ \exp(\beta(\phi_S(y) - \phi_f(y))) - \exp(\beta(V_{bs} - \phi_f(y))) \right\} \right]^{-\frac{1}{2}}$$

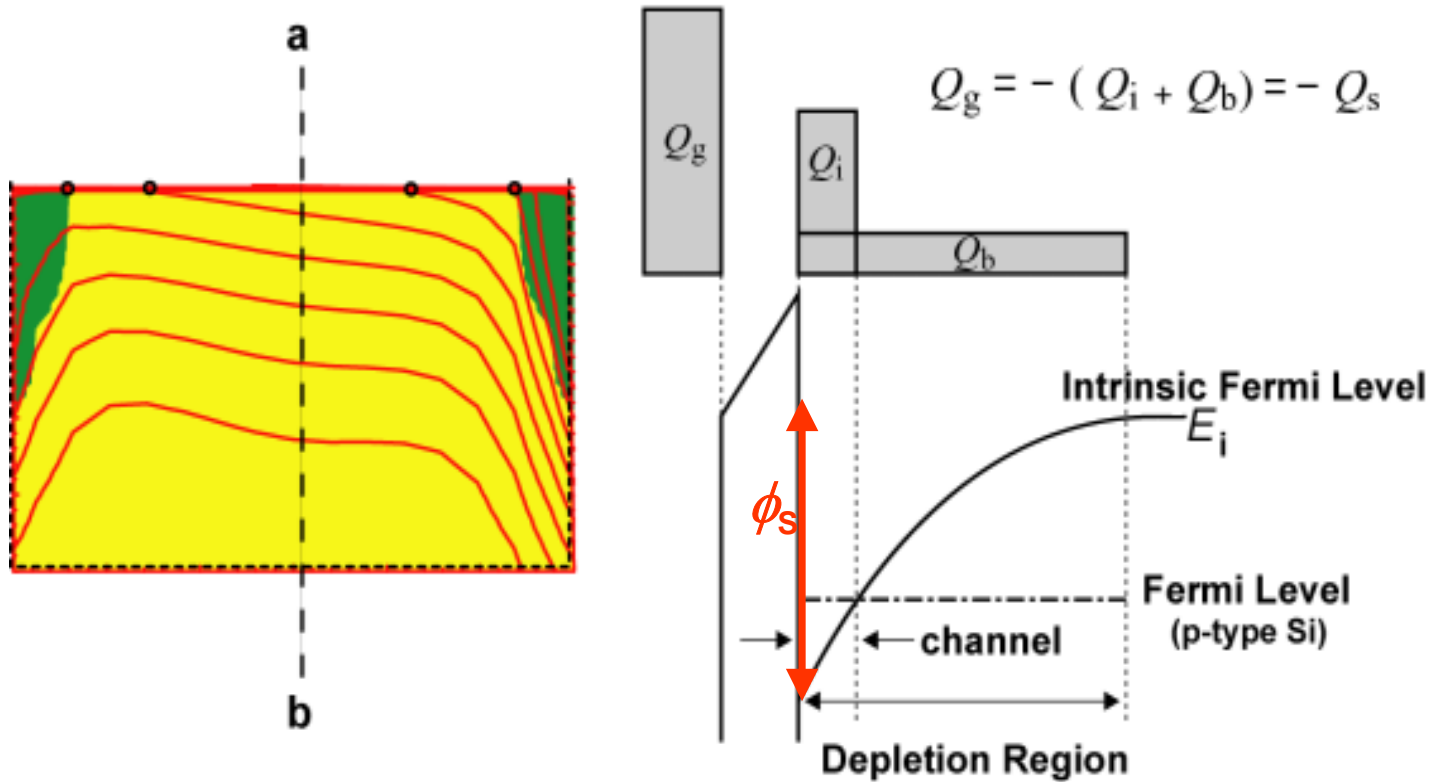
$$V'_G = V_{gs} - V_{fb}$$

$$\beta = \frac{q}{kT}$$



C. G. B. Garrett and W. H. Brattain, Phys. Rev., vol. 99, p. 376, 1955.

基板に誘起される全電荷



バルク電荷密度: $Q_b = \text{sqrt}(\beta\phi_s)$

反転層電荷密度: $Q_i = \text{function}(\phi_s)$ charge-sheet近似

デバイスの基本方程式

-Poisson:
$$\nabla^2 \phi = -\frac{q}{\epsilon_S} (N_D - N_A + p - n)$$

$$n = n_i \exp \frac{q(\phi - \phi_n)}{kT}$$

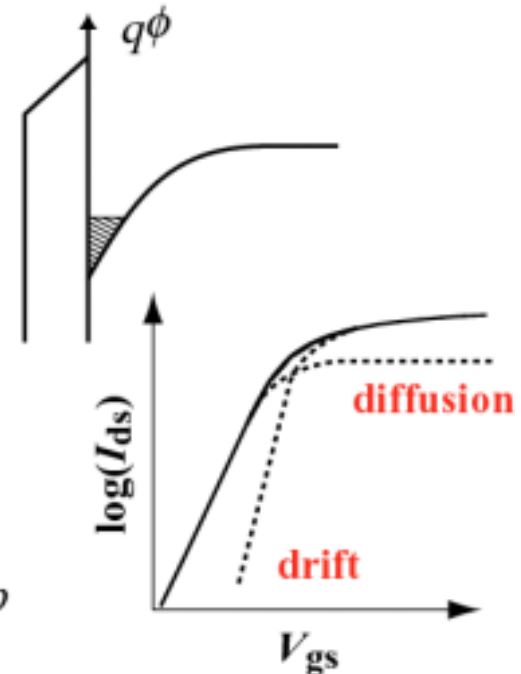
$$p = n_i \exp \frac{q(\phi_p - \phi)}{kT}$$

-Current Density:
$$j_n = -q \mu_n n \frac{\phi}{y} + q D_n \nabla n$$

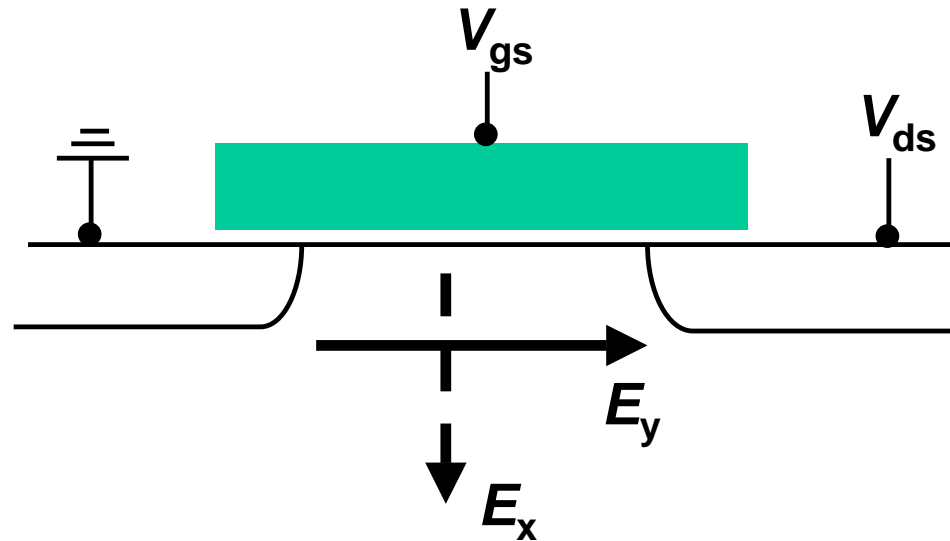
$$j_p = -q \mu_p p \frac{\phi}{y} - q D_p \nabla p$$

-Continuity:
$$\frac{\partial n}{\partial t} = G_n - R_n + \frac{1}{q} \nabla j_n = 0$$

$$\frac{\partial p}{\partial t} = G_p - R_p - \frac{1}{q} \nabla j_p = 0$$



ドリフト近似



電流: $I = qnv$

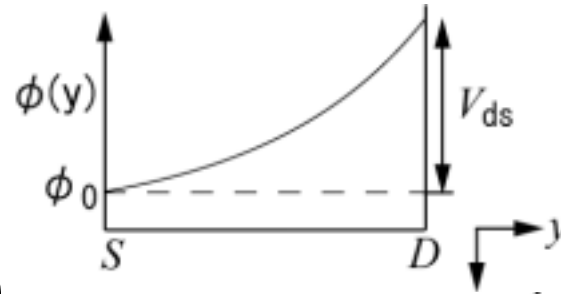
q : 電荷
 n : キャリア濃度
 v : キャリア速度

V_{gs} : キャリアの濃度を定める
 V_{ds} : キャリアの速さを定める

$$I_{ds} = W \mu q n \frac{d\phi}{dy}$$

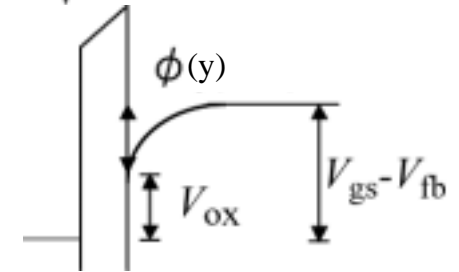
$$I_{ds} dy = W \mu q n d\phi$$

$$\phi(y) = \phi_0 + V(y)$$



Gaussの法則: $E_{ox} \epsilon_{ox} = E_{Si} \epsilon_{Si} = -Q_s$

$$\frac{V_{gs} - V_{fb} - \phi(y)}{T_{ox}} \epsilon_{ox} = -Q_n - Q_b$$



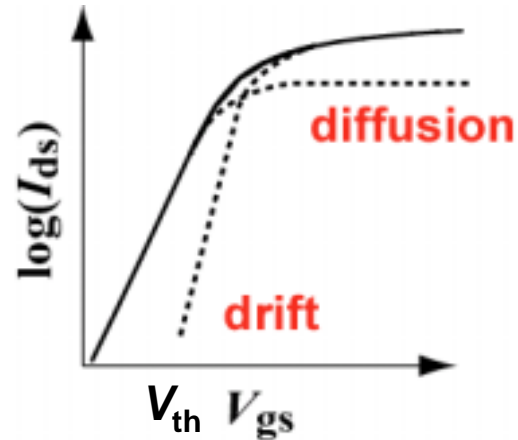
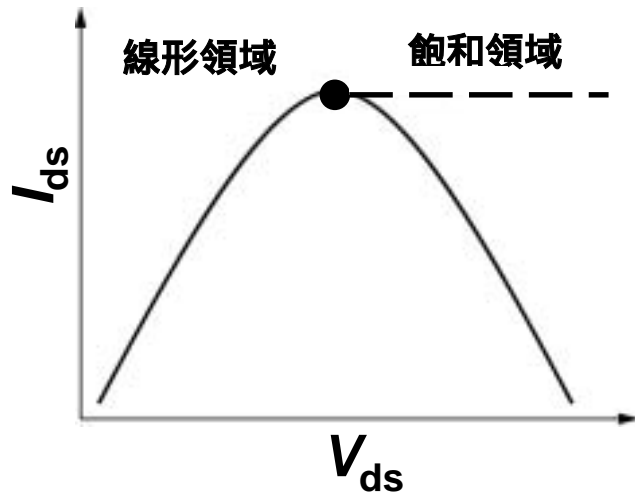
$$Q_n = qn = -\left[V_{gs} - V_{fb} - \phi_0 - V(y)\right] C_{ox} + \sqrt{2\epsilon_{Si} q N_{sub} \left[\phi_0 + V(y)\right]}$$

$$\int_0^L I_{ds} dy = W \mu \int_0^{V_{ds}} Q_n dV \Rightarrow I_{ds} = \mu \frac{W}{L} C_{ox} \left[(V_{gs} - V_{th}) V_{ds} - \frac{1}{2} V_{ds}^2 \right]$$

$$\phi_0 = 2\Phi_B; V_{th} = 2\Phi_B + \frac{\sqrt{2E_{Si} q N_{sub} 2\Phi_B}}{C_{ox}}$$

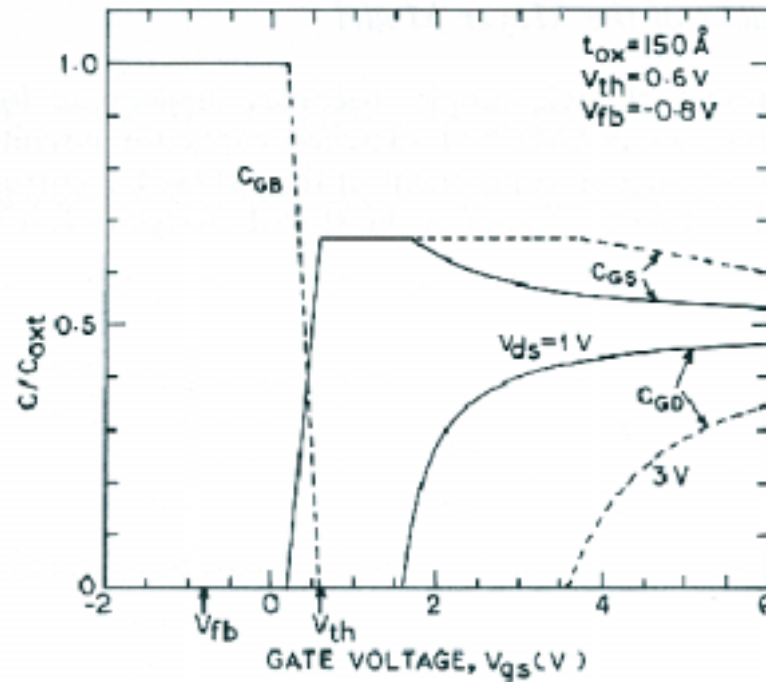
ドリフト近似の問題点

$$I_{ds} = \mu \frac{W}{L} C_{ox} \left[(V_{gs} - V_{th}) V_{ds} - \frac{1}{2} V_{ds}^2 \right]$$



キャパシタンスにおける問題

$$C_{gs} = \frac{2}{3} WLC_{ox} \left[1 - \frac{(V_{gs} - V_{th})^2}{(V_{gd} + V_{gs} - 2V_{th})^2} \right]$$



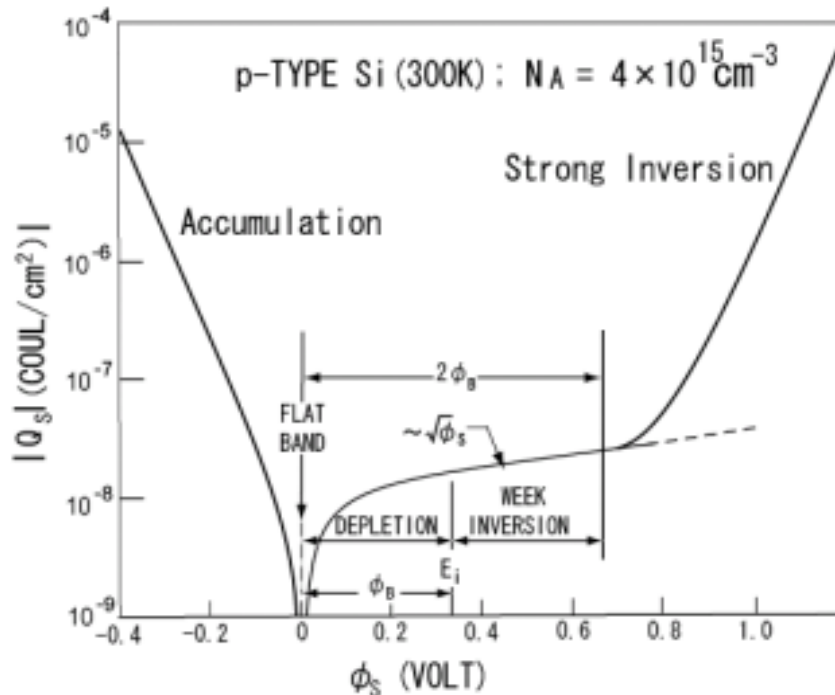
MOSキャパシタンス

$$-Q_S = C_{ox}(V'_G - \phi_S(y)) = \sqrt{\frac{2\epsilon_s q N_{sub}}{\beta}} \left[\exp\{-\beta(\phi_S(y) - V_{bs})\} + \beta(\phi_S(y) - V_{bs}) - 1 \right. \\ \left. + \frac{n_{p0}}{p_{p0}} \left\{ \exp(\beta(\phi_S(y) - \phi_f(y))) - \exp(\beta(V_{bs} - \phi_f(y))) \right\} \right]^{-\frac{1}{2}}$$

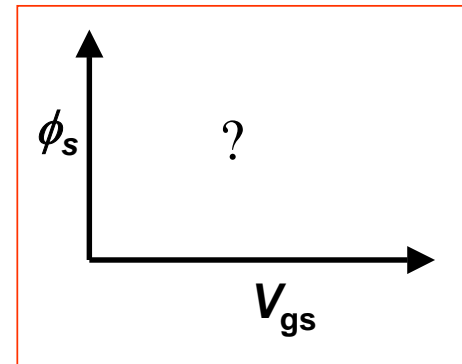
$$V'_G = V_{gs} - V_{fb}$$

$$\beta = \frac{q}{kT}$$

ϕ_f : 擬フェルミ(0 V_{ds})

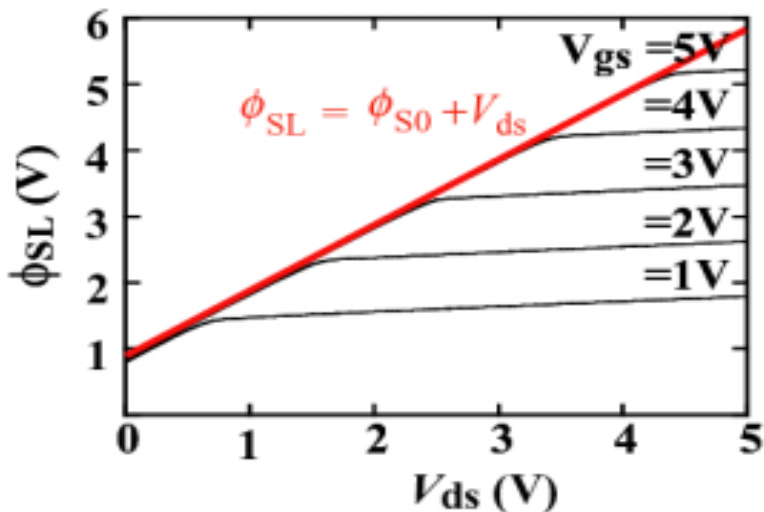
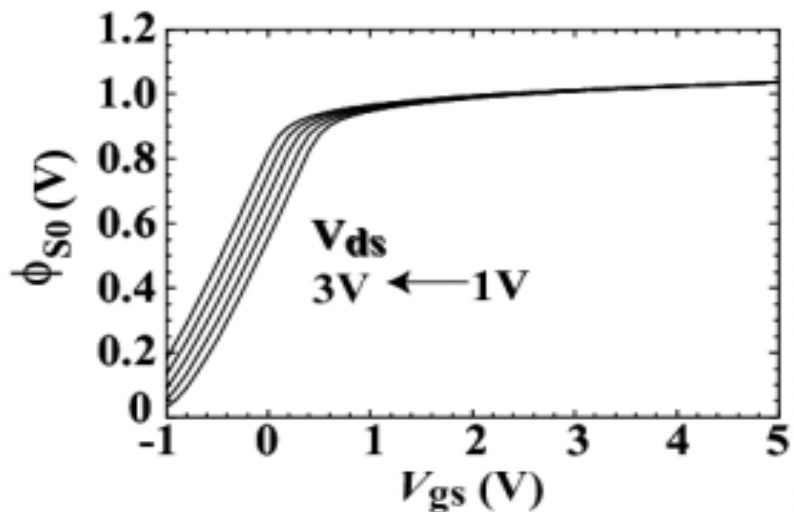
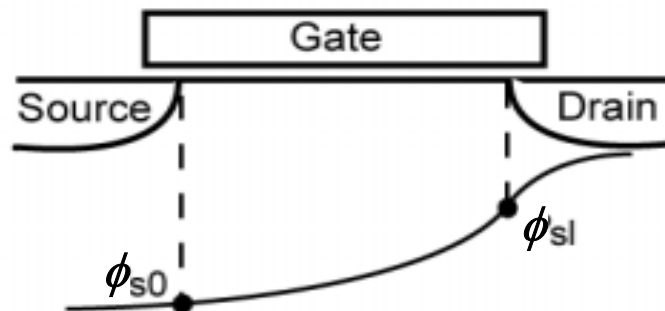
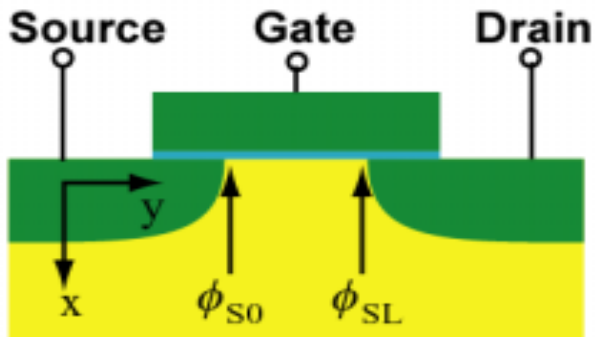


C. G. B. Garrett and W. H. Brattain, Phys. Rev., vol. 99, p. 376, 1955.



反復計算によって解く

表面ポテンシャルの計算結果



ϕ_{sl} : Gradual-Channel近似的の終わり

ドリフト・拡散モデルの電流式導出

$$I_{\text{ds}}(y) = W_{\text{eff}} q \mu n(y) \left(-\frac{d\phi_s(y)}{dy} + \frac{1}{\beta} \frac{d \ln n(y)}{dy} \right)$$

$$I_{\text{ds}} = \int_0^L I_{\text{ds}}(y) dy$$

$$\left(\begin{array}{l} \frac{d \ln n(y)}{dy} = \frac{1}{n(y)} \frac{dn(y)}{dy} \\ \frac{I_{\text{ds}}}{W_{\text{eff}} q \mu} + n(y) \beta \frac{d\phi_s(y)}{dy} = \frac{dn(y)}{dy} \\ -C_{\text{ox}} [V_G - \phi_s(y)] = Q_B + qn(y) \end{array} \right)$$

電流はポテンシャルの関数

ドレイン電流式

$$I_{ds} = W_{\text{eff}} q \mu n(y) \left(-\frac{d\phi_S(y)}{dy} + \frac{1}{\beta} \frac{d \ln n(y)}{dy} \right)$$

$$I_{ds} = \frac{W_{\text{eff}}}{L_{\text{eff}}} \mu \frac{IDD}{\beta}$$

$$\begin{aligned} IDD &= C_{\text{ox}} (\beta V_G' + 1) (\phi_{\text{SL}} - \phi_{\text{S0}}) - \frac{\beta}{2} C_{\text{ox}} (\phi_{\text{SL}}^2 - \phi_{\text{S0}}^2) \\ &- \frac{2}{3} \sqrt{\frac{2\epsilon_s q N_{\text{sub}}}{\beta}} \left[\{\beta(\phi_{\text{SL}} - V_{\text{bs}}) - 1\}^{\frac{3}{2}} - \{\beta(\phi_{\text{S0}} - V_{\text{bs}}) - 1\}^{\frac{3}{2}} \right] \\ &+ \sqrt{\frac{2\epsilon_s q N_{\text{sub}}}{\beta}} \left[\{\beta(\phi_{\text{SL}} - V_{\text{bs}}) - 1\}^{\frac{1}{2}} - \{\beta(\phi_{\text{S0}} - V_{\text{bs}}) - 1\}^{\frac{1}{2}} \right] \end{aligned}$$

電流式は表面ポテンシャルの関数

J. R. Brews, Solid-State Electron. , vol. 21, p.345, 1978;

G. Bacarani, IEEE J. Solid- State Electron Devices, vol. 2, p.62, 1978;

F. van der Wiele, Solid-State Electron., vol.22, p.991, 1979.

ドレイン電流式の比較

Drift-Diffusion Approximation: HiSIM

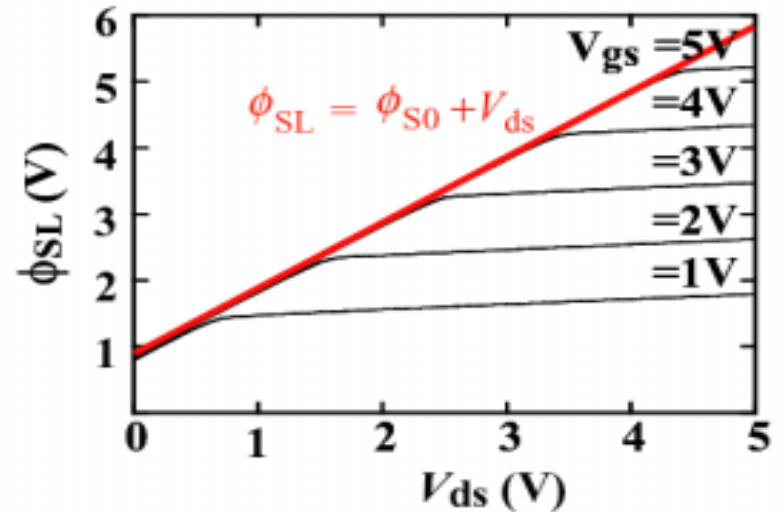
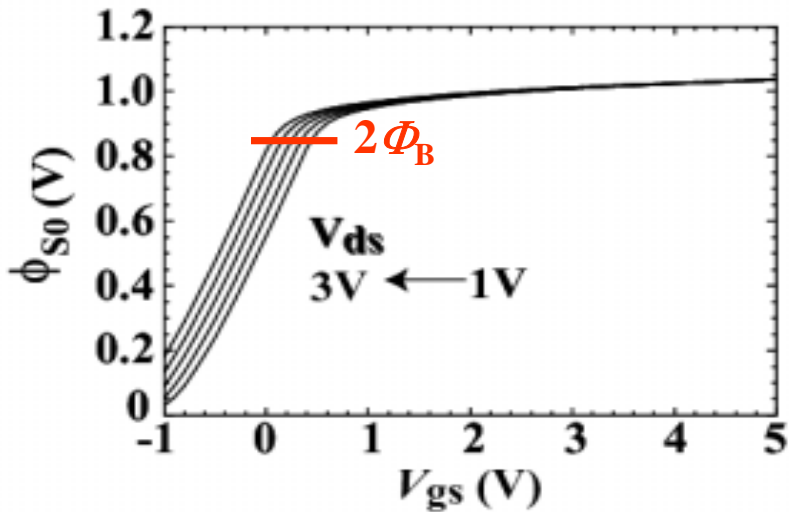
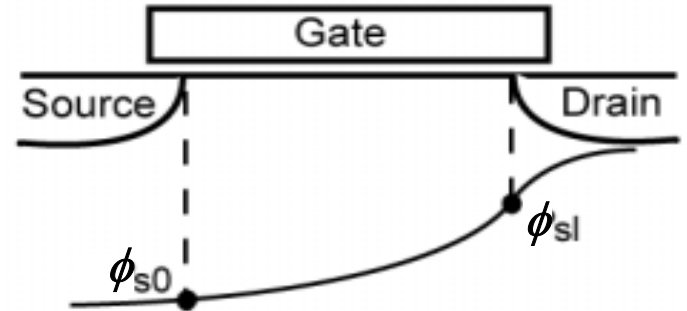
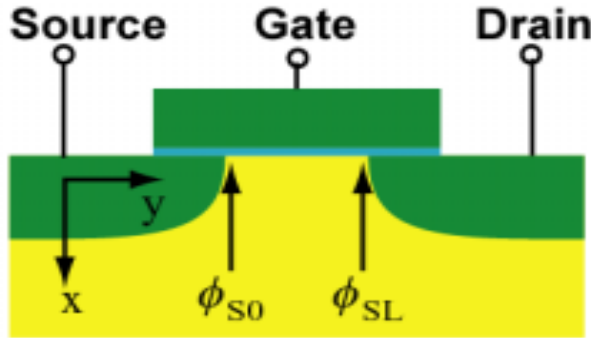
$$I_{ds} \left(\frac{1}{\beta} \mu \frac{W}{L} \right) = C_{ox}(1 + \beta V'_G)(\phi_{SL} - \phi_{S0}) - \frac{\beta}{2} C_{ox}(\phi_{SL}^2 - \phi_{S0}^2) \\ - \frac{2}{3} \sqrt{\frac{2\epsilon_s q N_{sub}}{\beta}} [(\beta \phi_{SL} - 1)^{\frac{3}{2}} - (\beta \phi_{S0} - 1)^{\frac{3}{2}}] \\ + \sqrt{\frac{2\epsilon_s q N_{sub}}{\beta}} [(\beta \phi_{SL} - 1)^{\frac{1}{2}} - (\beta \phi_{S0} - 1)^{\frac{1}{2}}] \\ V'_G = V_{gs} - V_{fb} + (\Delta V_{th})$$

Drift Approximation: $\phi_{S0} = 2\Phi_B = \frac{2}{\beta} \ln\left(\frac{N_{sub}}{n_i}\right)$; $\phi_{SL} = \phi_{S0} + V_{ds}$

$$I_{ds} \left(\mu \frac{W}{L} C_{ox} \right) = (V_G - 2\Phi_B - \frac{\sqrt{2\epsilon_s q N_{sub}}}{C_{ox}} 2\Phi_B^{\frac{1}{2}}) V_{ds} - \left(\frac{1}{2} + \frac{\sqrt{2\epsilon_s q N_{sub}}}{4C_{ox}} 2\Phi_B^{\frac{1}{2}} \right) V_{ds}^2 \\ \simeq (V_G - V_{th}) V_{ds} - \frac{1}{2} V_{ds}^2$$

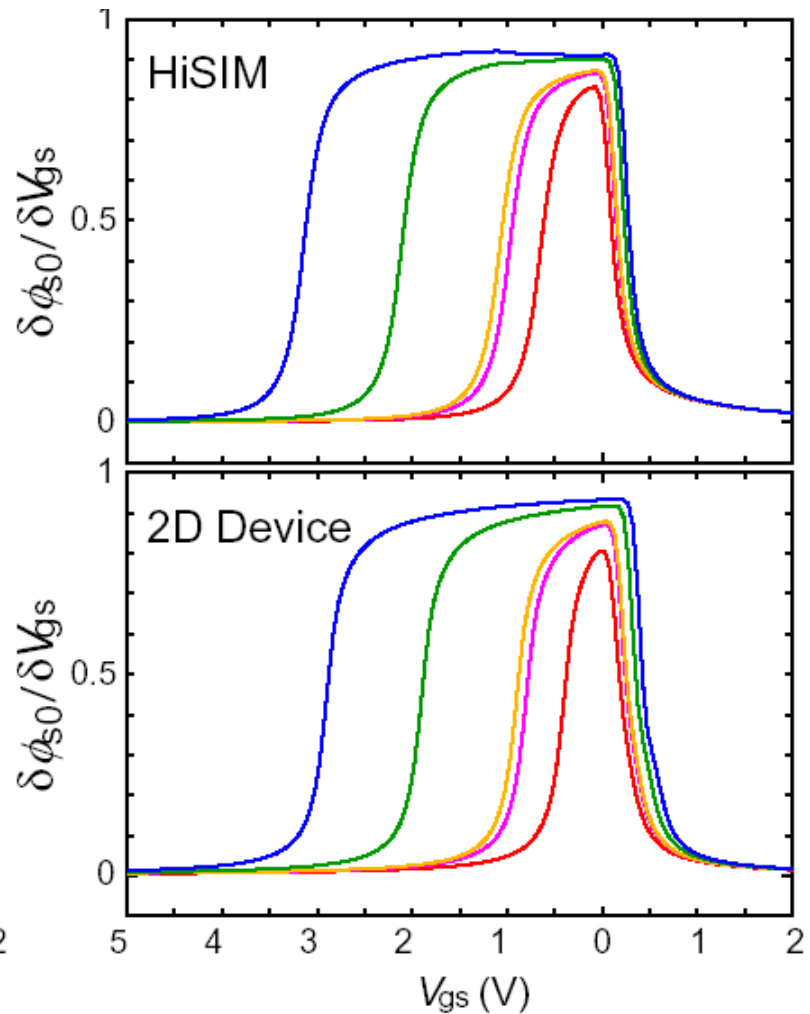
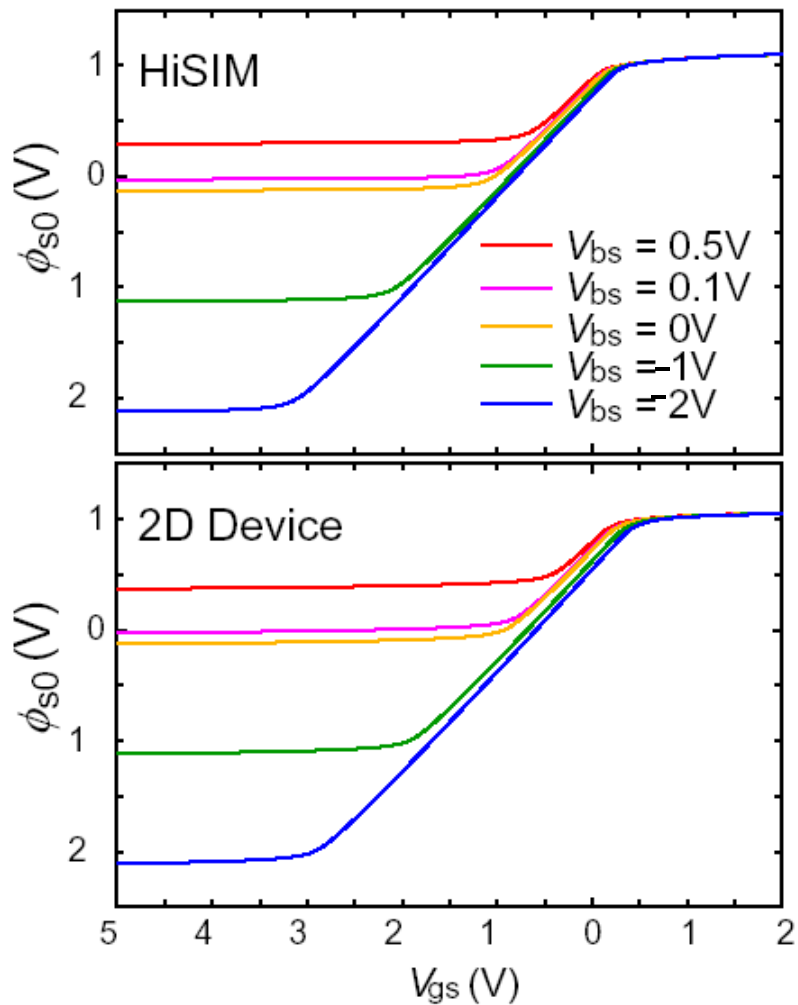
$$V_G = V_{gs} - V_{fb}$$

ドリフト近似の表面ポテンシャル

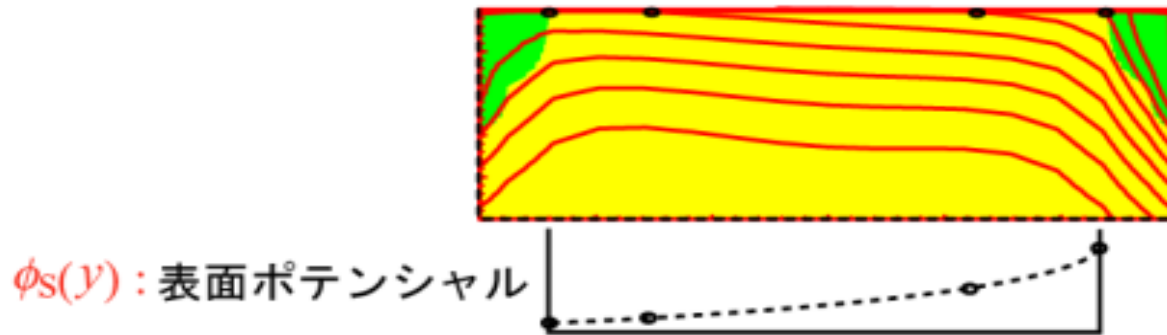


ϕ_{SL} : Gradual-Channel近似的の終わり

表面ポテンシャルとその微分



チャンネル方向のポテンシャル分布



Charge-Sheet近似

Poisson's Equation:

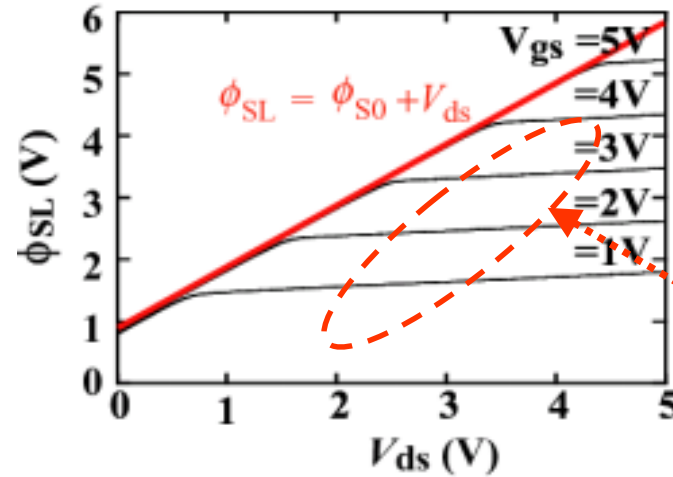
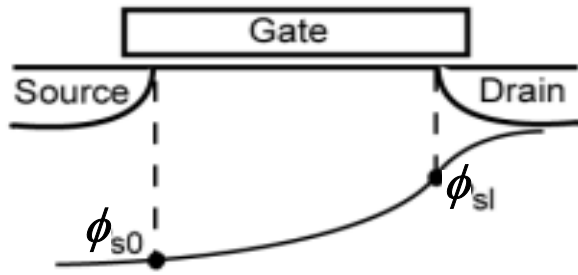
$$Q_b(y) = -qN_{\text{sub}} \times W_d = -\sqrt{\frac{2\epsilon_s qN_{\text{sub}}}{\beta}} \left\{ \beta (\phi_s(y) - V_{\text{bs}}) - 1 \right\}^{-\frac{1}{2}}$$

Gauss's Law:

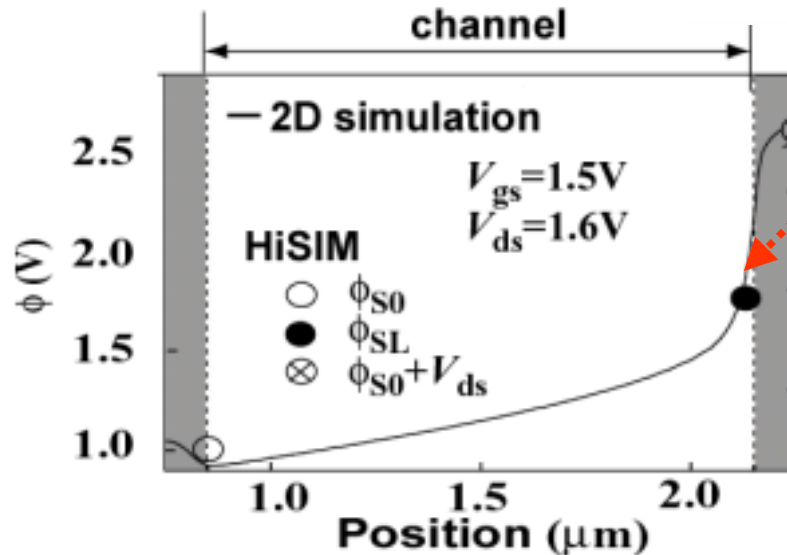
$$Q_g(y) = C_{\text{ox}} (V_G' - \phi_s(y))$$

$$Q_g(y) = -(Q_i(y) + Q_b(y))$$

Gradual-Channel近似



飽和領域



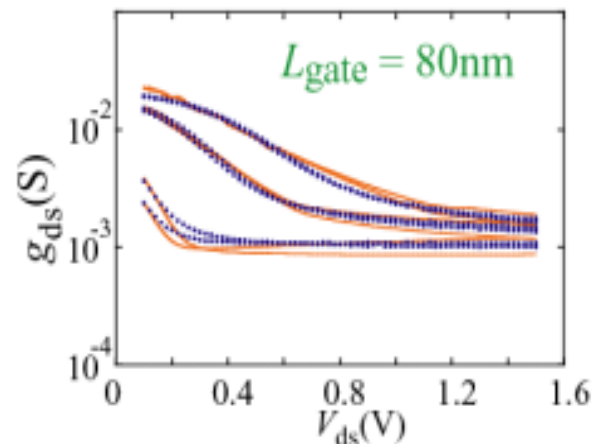
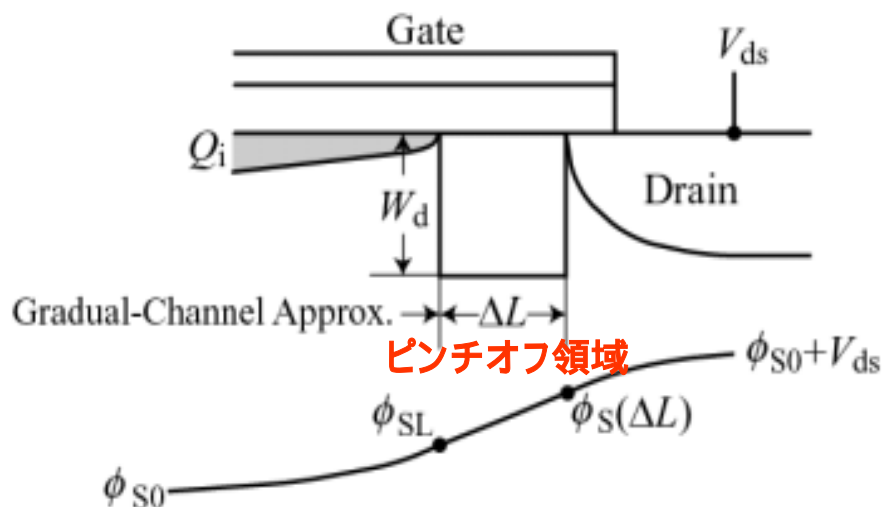
ピンチオフ領域

ϕ_{SL} : End of Gradual-Channel Approx.

チャンネル内位置?

Channel-Length Modulation

Gradual-Channel 近似の拡張



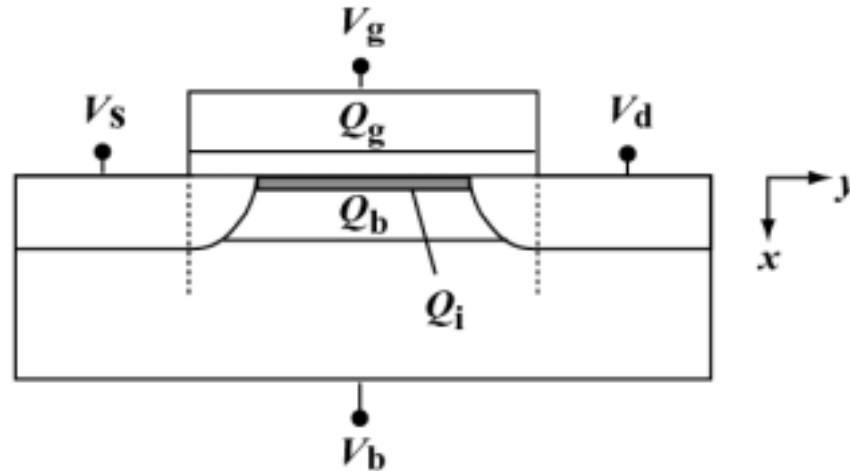
ポテンシャル値と位置を同時に決定できない。

ポテンシャルを定義： $\phi_{S(\Delta L)} = CLMI (\phi_{S0} + V_{ds}) + (1 - CLMI) \phi_{SL}$

($CLMI$ はジャンクション形状依存)

$$\rightarrow \Delta L = \epsilon_S \frac{E(\Delta L) - E(0)}{qN_{sub} + Q_i/W_d}$$

ノードに誘起される電荷



$$Q_I = W \int_0^L Q_i(y) dy;$$

$$Q_B = W \int_0^L Q_b(y) dy$$

$$Q_S = W \int_0^L \left(1 - \frac{y}{L}\right) Q_i(y) dy;$$

$$Q_D = W \int_0^L \frac{y}{L} Q_i(y) dy$$

$$I_{ds} = q \frac{W}{L} \int v Q_i dy$$

current



$$Q$$

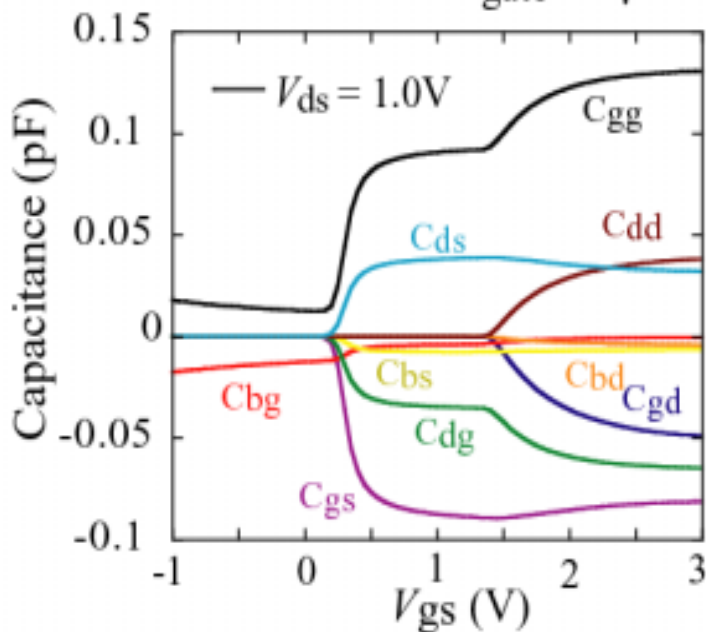
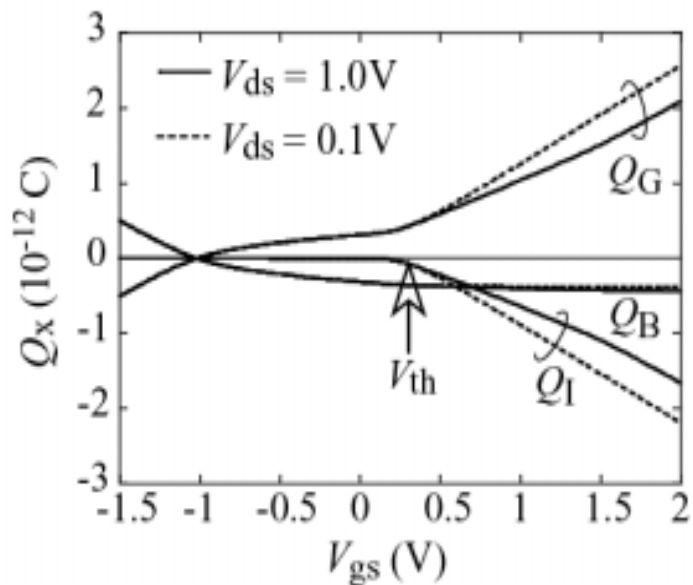
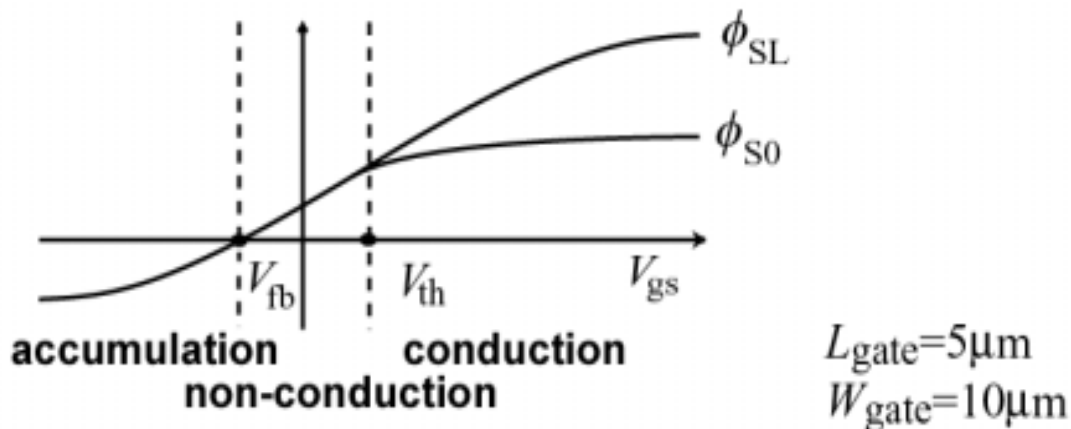
charge



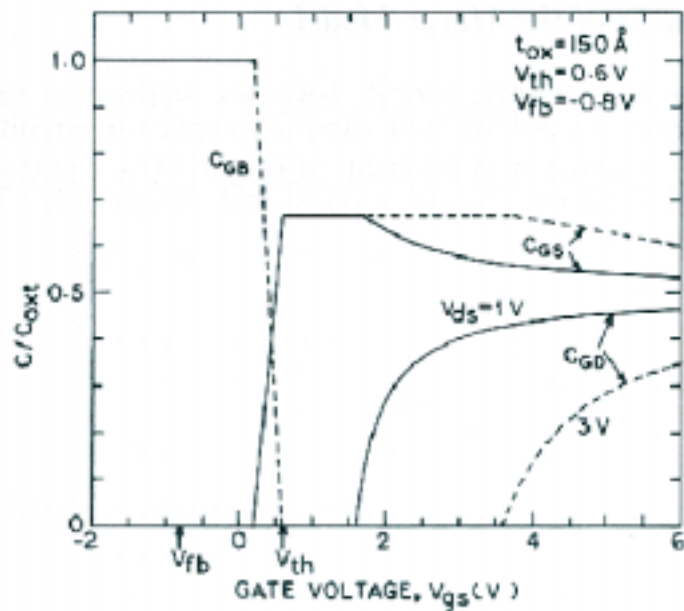
$$C_{jk} = \frac{dQ_j}{dV_k}$$

capacitance

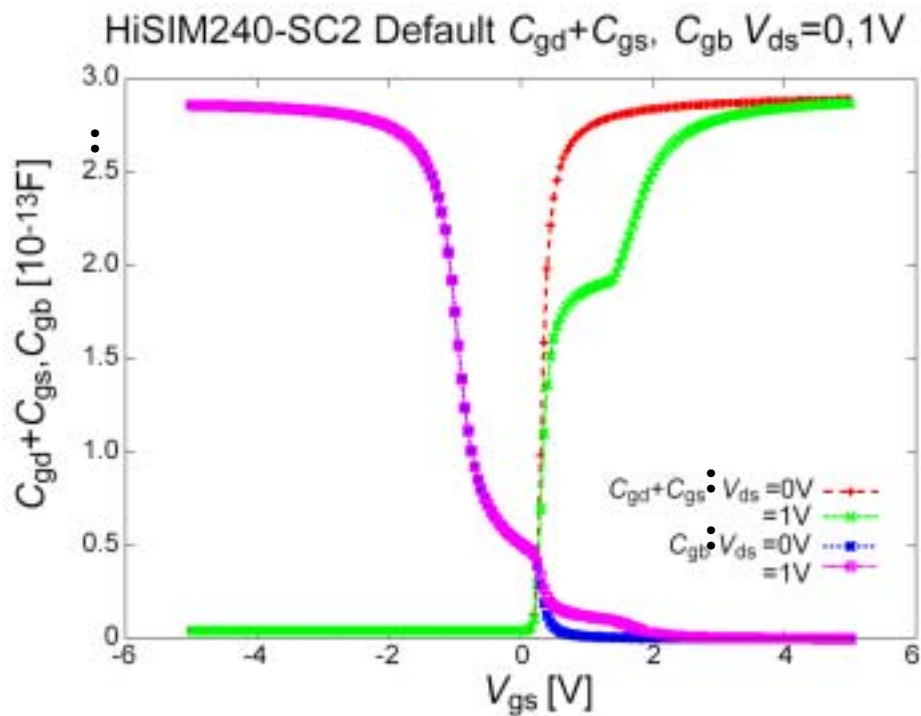
電荷とキャパシタンス



Driftモデル

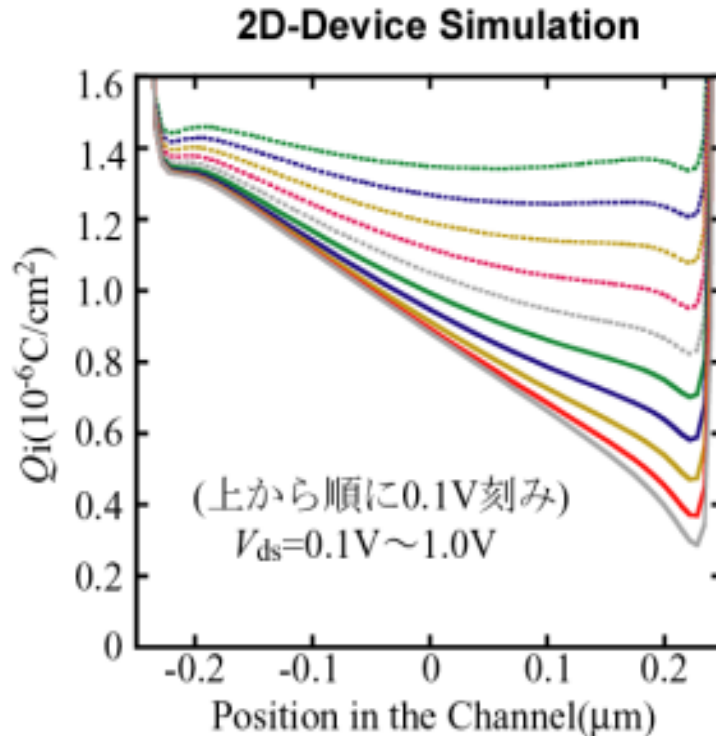


表面ポテンシャルモデル



電荷の分配

$$I_x(t) = I_x(V(t)) + dQ_x/dt$$



$$Q_I = Q_S + Q_D$$

$$Q_S = W \int_0^L \left(1 - \frac{y}{L}\right) Q_i(y) dy$$

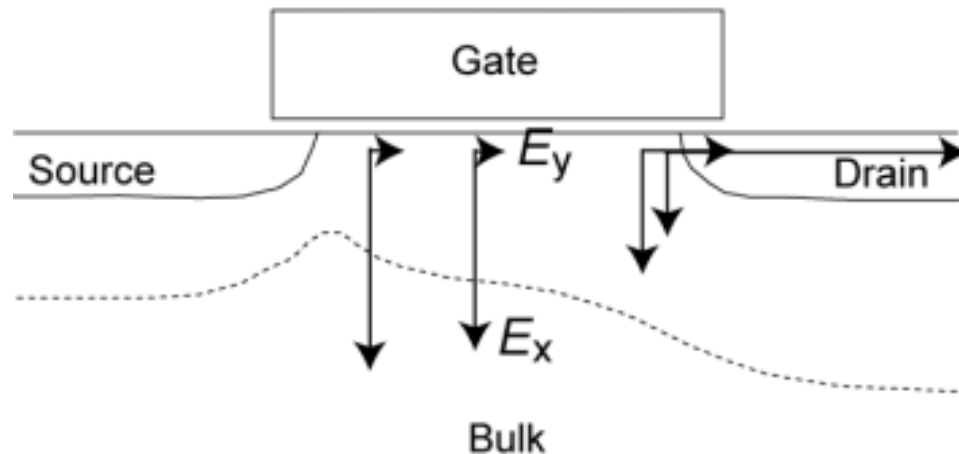
$$Q_D = W \int_0^L \frac{y}{L} Q_i(y) dy$$

S. - Y. Oh, D. E. Ward, and R.W. Dutton,
IEEE J. Solid-State Circuits, vol. SC-15, p.636, 1980.

移動度モデル

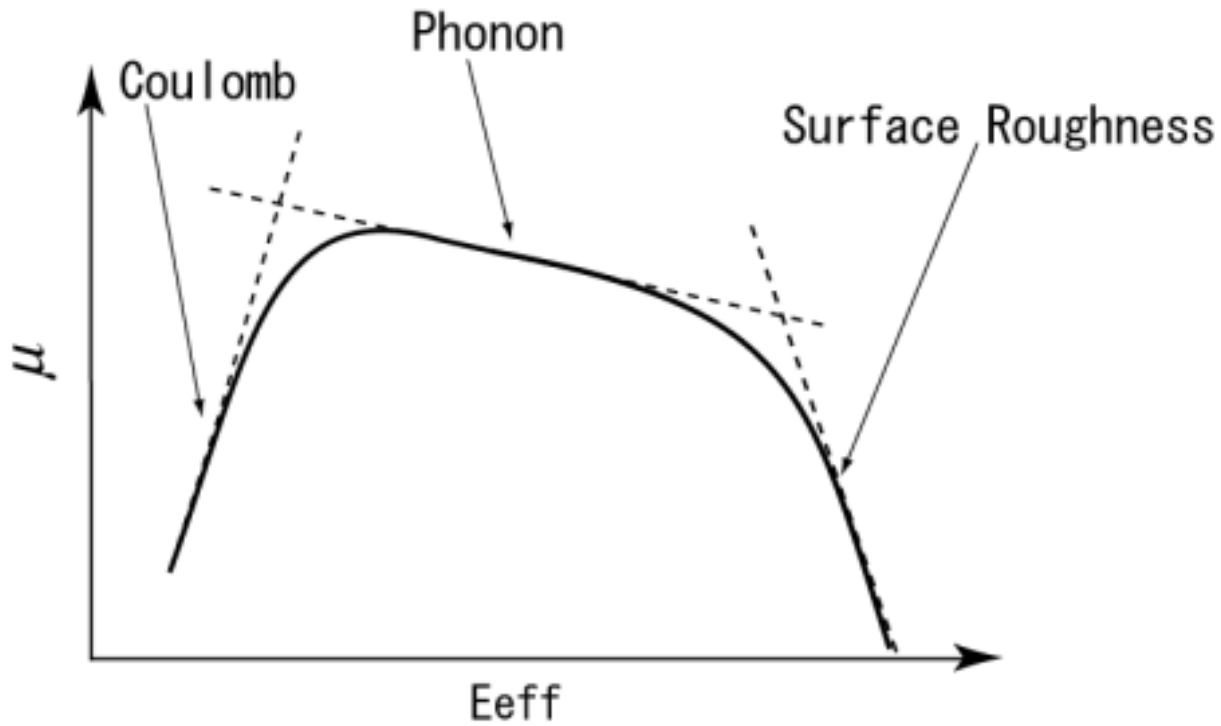
$$I_{ds} = \mu \frac{W}{L} C_{ox} \left[(V_{gs} - V_{th}) V_{ds} - \frac{1}{2} V_{ds}^2 \right]$$

μ : mobility (cm²/Vs)



- low-field mobility
- high-field mobility

低電界移動度を支配するキャリア散乱



$$\frac{1}{\mu_0} = \frac{1}{\mu_{CB}} + \frac{1}{\mu_{PH}} + \frac{1}{\mu_{SR}} \quad : \text{Matthiessen Rule}$$

$$\bullet \mu_{CB} \text{ (Coulomb)} = MUECB0 + MUECBI \frac{Q_i}{q \times 10^{11}}$$

$$\bullet \mu_{PH} \text{ (phonon)} = \frac{MUEPH0}{(T/300K)^{MUETMP} \times E^{eff} MUEPHI}$$

$$\bullet \mu_{SR} \text{ (surface roughness)} = \frac{MUESR0}{E_{eff} MUESRI}$$

$$E_{eff} = \frac{1}{\epsilon_S} (NDEP \times Q_b + NINV \times Q_i)$$

$$NINV = NINV - NINVD \times V_{ds}$$

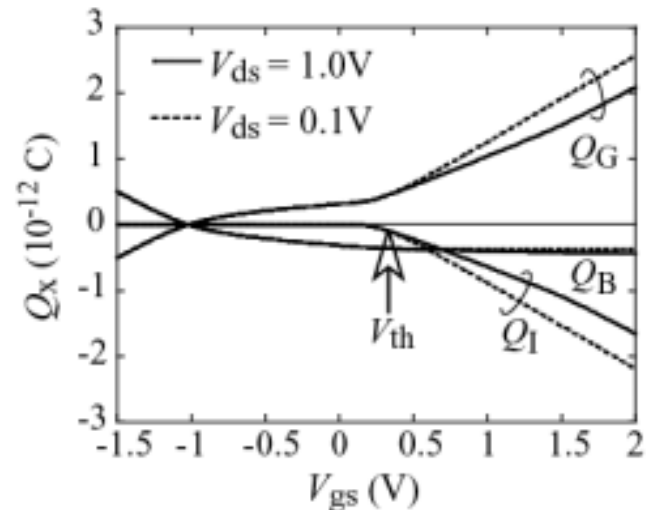
Universality: $MUEPHI = 0.3$

$MUESRI = 2.0$

+

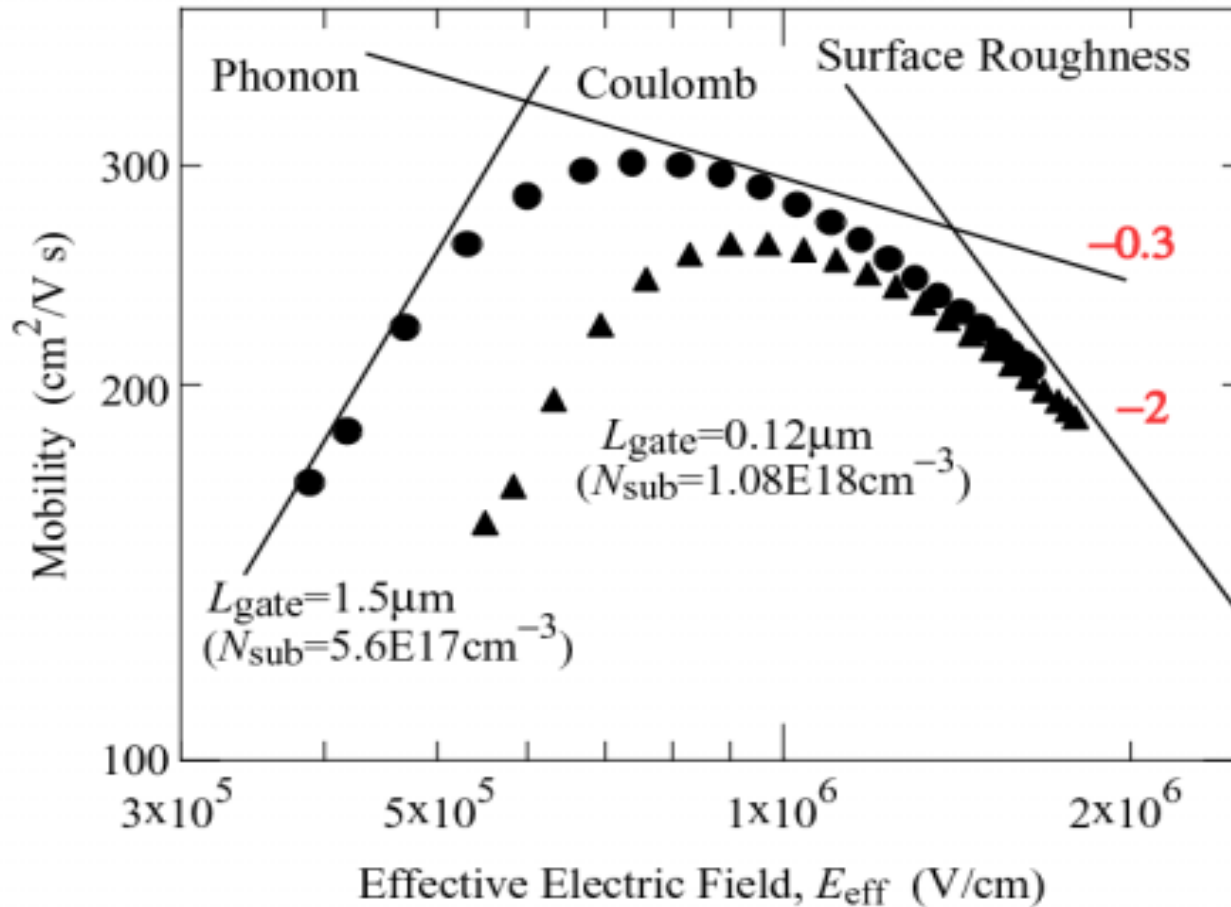
$NDEP = 1.0$

$NINV = 0.5$



E_{eff} はほぼ V_{gs} の線形関数

実測値の考察



高電界移動度

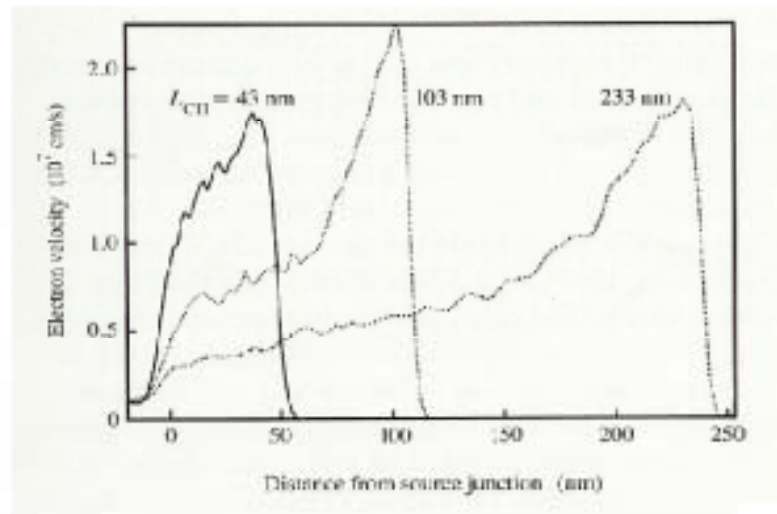
$$\mu = \frac{\mu_0}{\left(1 + \left(\frac{\mu_0 E_y}{V_{MAX}}\right)^{BB}\right)^{\frac{1}{BB}}} \quad \begin{array}{l} BB = 2.0 \text{ (electron)} \\ 1.0 \text{ (hole)} \end{array}$$

$$V_{MAX} = \frac{V_{MAX}}{(1.8 + 0.4 \times (T/300K) + 0.1 \times (T/300K)^2)}$$

:Caughtey-Thomas

$$V_{MAX} = V_{MAX} (1.0 + V_{OVER} / L_{gate} V_{OVERP})$$

Velocity Overshoot



S. E. Laux et al., IBM J. R&D, vol.34, p. 466, 1990.

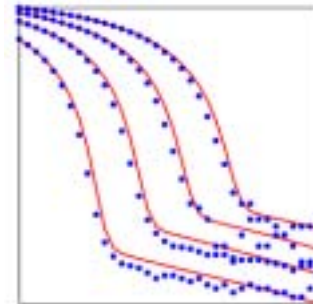
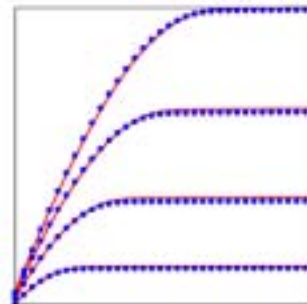
実測電流値と計算値の比較

NMOS
Large

— HiSIM2
○ Measurement

$I_d V_d$

gds



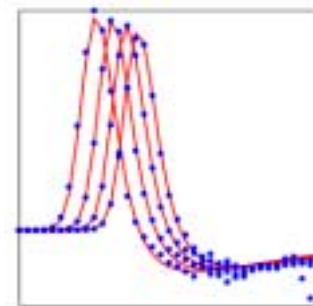
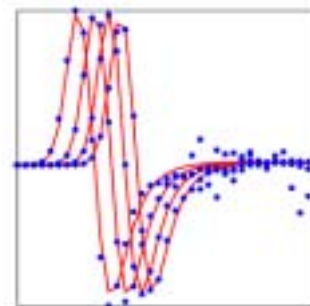
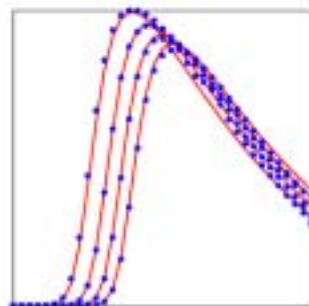
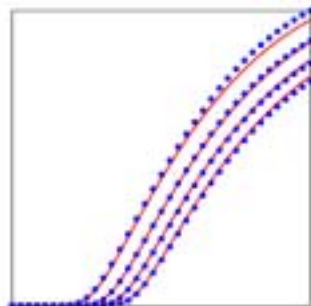
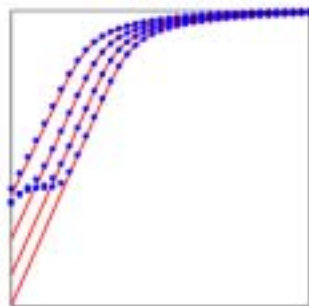
$I_d V_g$ _low

$I_d V_g$ _low

gm_low

gm2_low

gm3_low



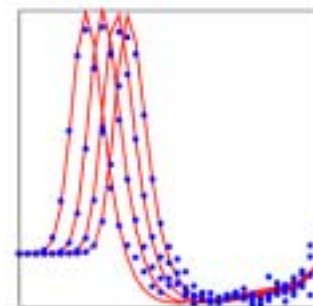
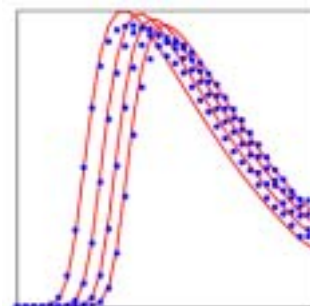
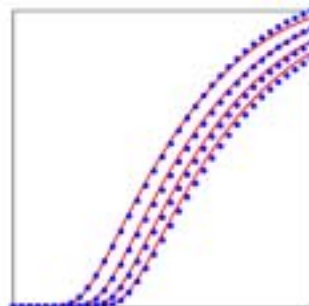
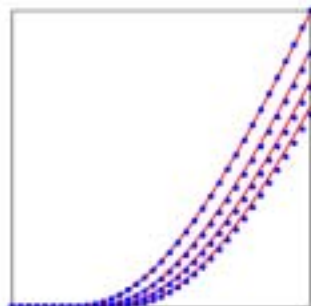
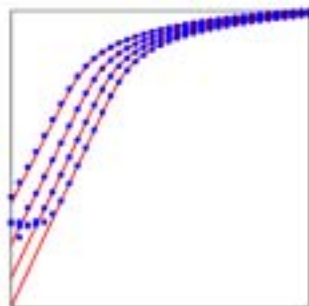
$I_d V_g$ _high

$I_d V_g$ _high

gm_high

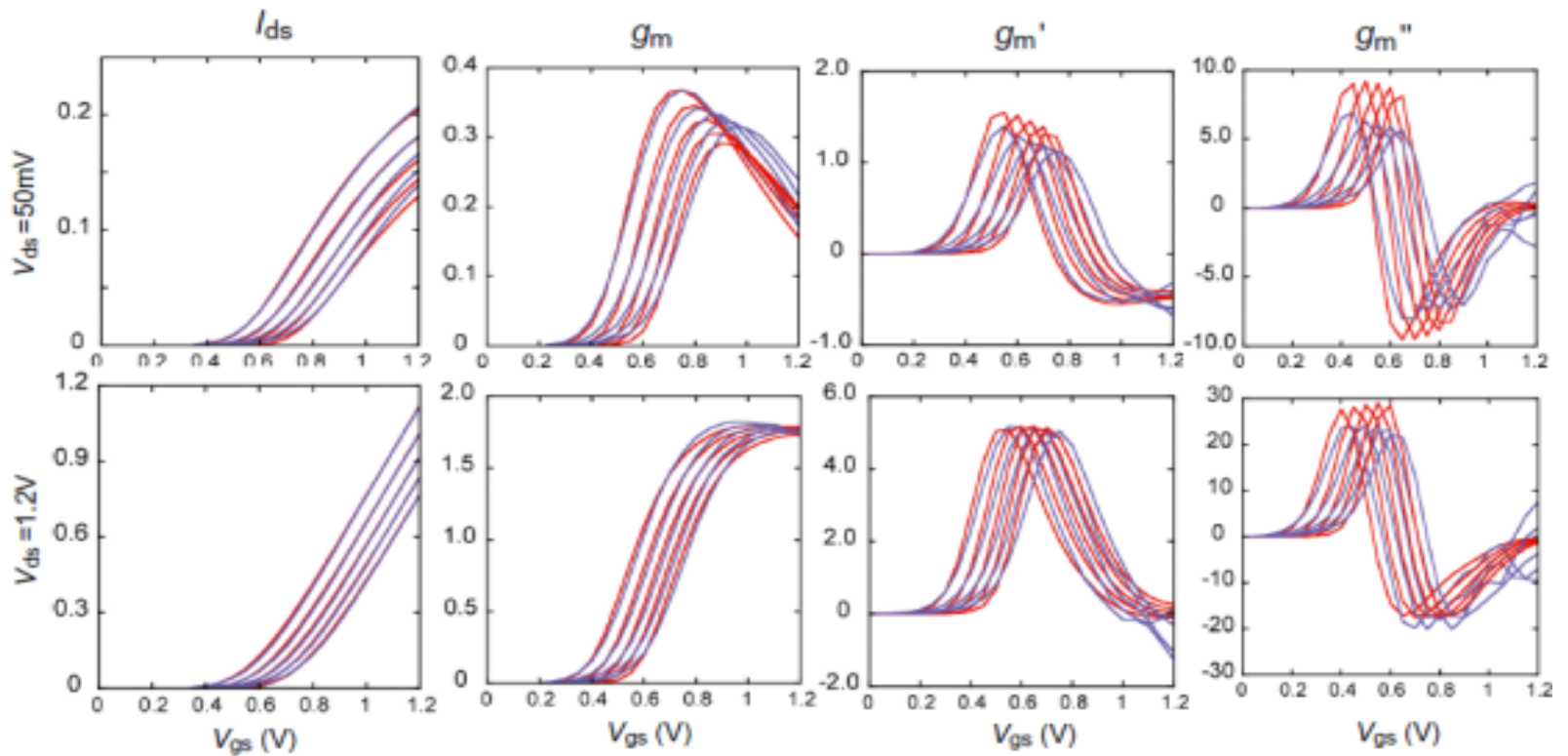
gm2_high

gm3_high

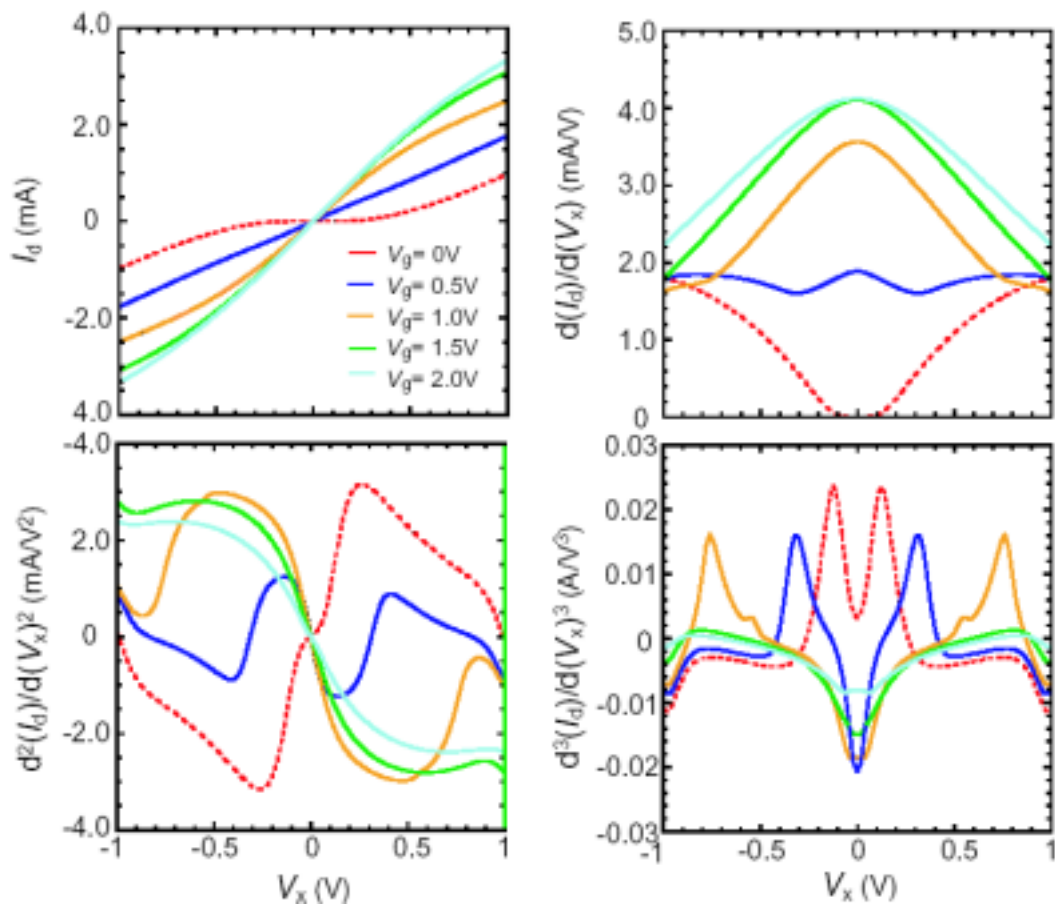


微分特性

$L_{\text{gate}}=40\text{nm}$



対称性



ドリフト近似では実現できない

$$I_{ds} = \mu \frac{W}{L} C_{ox} \left[(V_{gs} - V_{th}) V_{ds} - \frac{1}{2} V_{ds}^2 \right]$$

基本方程式

-Poisson:
$$\nabla^2 \phi = -\frac{q}{\epsilon_{Si}} (N_D - N_A + p - n)$$

$$n = n_i \exp \frac{q(\phi - \phi_n)}{kT}$$

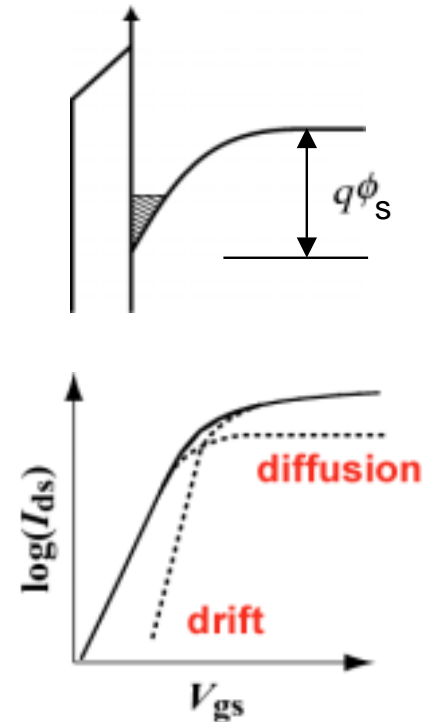
$$p = n_i \exp \frac{q(\phi_p - \phi)}{kT}$$

-Current Density:
$$j_n = q\mu_n n \frac{\phi}{y} + qD_n \nabla n$$

$$j_p = q\mu_p p \frac{\phi}{y} - qD_p \nabla p$$

-Continuity:
$$I(t) = I_0(t) + \frac{dQ}{dt}$$

(solved by SPICE)



すべてのデバイス特性はポテンシャルの関数

Donald O. Pederson



**You don't get any credit
for doing 95 percent of the job.**

SPICE: Simulation Program with Integrated Circuit Emphasis

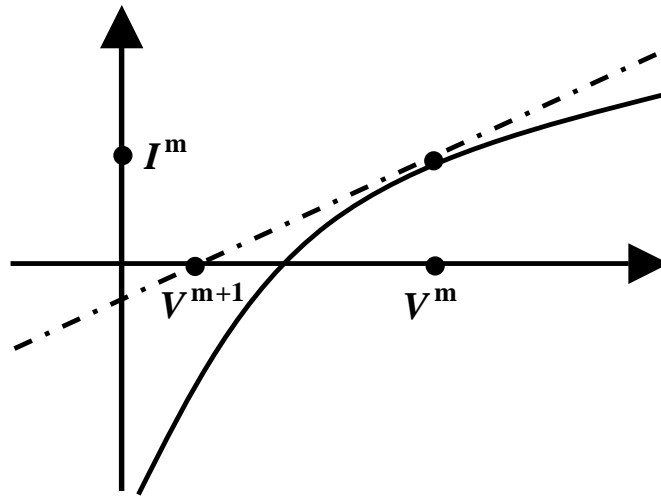
SPICEシミュレーションの基礎

Circuit Simulator :

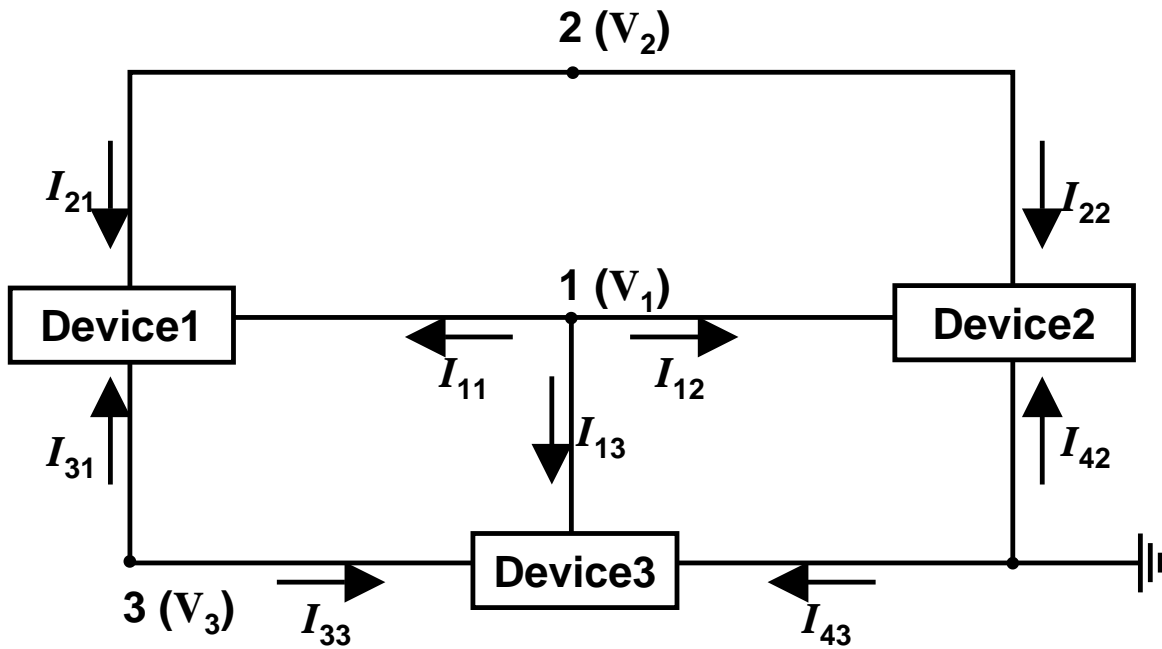
Finding a set of $V'_s = (V_1, V_2, V_3, \dots, V_n)$

for $I'_s(I_1, I_2, I_3, \dots, I_n)$ by fulfilling the kirchhof law.

$$V^{m+1} = V^m - \frac{I^m}{g^m}$$
$$g^m V^{m+1} = -I^m + g^m V^m$$



Newton's Method



$$\begin{aligned}
 I_1 &= I_{11} + I_{12} + I_{13} \\
 I_2 &= I_{21} + I_{22} \\
 I_3 &= I_{31} + I_{13}
 \end{aligned}$$

$$\begin{bmatrix}
 +\frac{\partial I_{11}^m}{\partial V_1} + \frac{\partial I_{11}^m}{\partial V_2} + \frac{\partial I_{11}^m}{\partial V_3} \\
 +\frac{\partial I_{21}^m}{\partial V_1} + \frac{\partial I_{21}^m}{\partial V_2} + \frac{\partial I_{21}^m}{\partial V_3} \\
 +\frac{\partial I_{31}^m}{\partial V_1} + \frac{\partial I_{31}^m}{\partial V_2} + \frac{\partial I_{31}^m}{\partial V_3}
 \end{bmatrix}
 \begin{bmatrix}
 V_1^{m+1} \\
 V_2^{m+1} \\
 V_3^{m+1}
 \end{bmatrix}
 = -
 \begin{bmatrix}
 +I_{11}^m \\
 +I_{21}^m \\
 +I_{31}^m
 \end{bmatrix}
 +
 \begin{bmatrix}
 +\frac{\partial I_{11}^m}{\partial V_1} + \frac{\partial I_{11}^m}{\partial V_2} + \frac{\partial I_{11}^m}{\partial V_3} \\
 +\frac{\partial I_{21}^m}{\partial V_1} + \frac{\partial I_{21}^m}{\partial V_2} + \frac{\partial I_{21}^m}{\partial V_3} \\
 +\frac{\partial I_{31}^m}{\partial V_1} + \frac{\partial I_{31}^m}{\partial V_2} + \frac{\partial I_{31}^m}{\partial V_3}
 \end{bmatrix}
 \begin{bmatrix}
 V_1^m \\
 V_2^m \\
 V_3^m
 \end{bmatrix}$$